

Forward Electron Spectroscopy at the ESR

- **J.R.Macdonald Laboratory, Kansas State University**
S. Hagmann , P. Richard, C.L. Cocke,
- **GSI- Darmstadt**
Th. Stoehlker, R. Mann, Ch. Kozhuharov
- **Max Planck Inst. f. Kernphysik-Heidelberg**
H. Kollmus, J. Ullrich, R. Moshhammer
- **Inst.f. Kernphysik, Universität Frankfurt**
H. Schmidt-Böcking , S.H.

Overview

A. Motivation:

Kinematically complete Ionization Experiments

many-electron continua ($2 e^- + 1 n$)

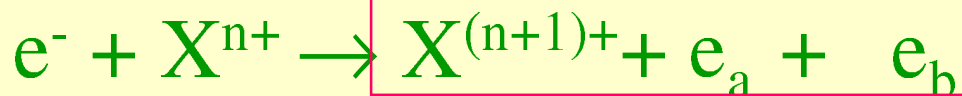
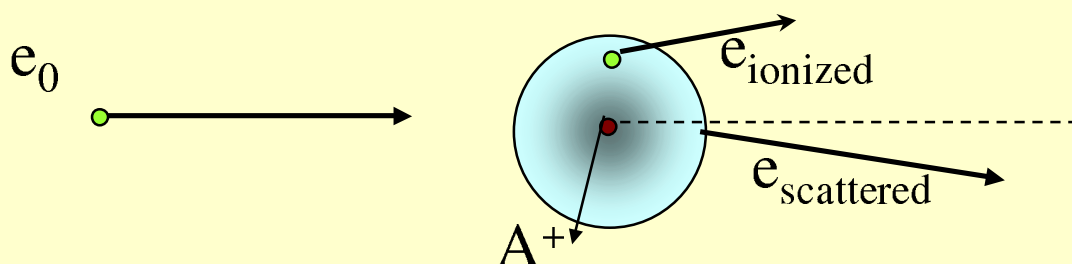
B. A Tool :

Zero Degree Spectroscopy,

from autoionizing Rydberg states to Quasi-Photoionization of Fe XVII

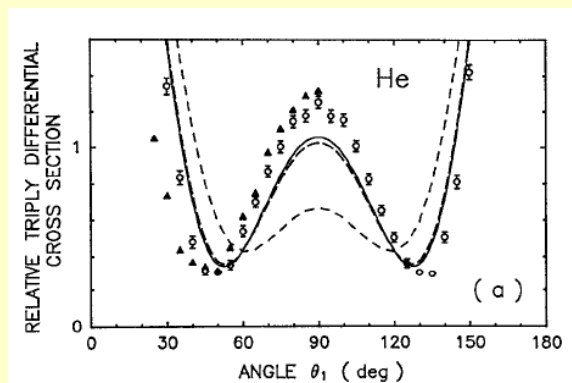
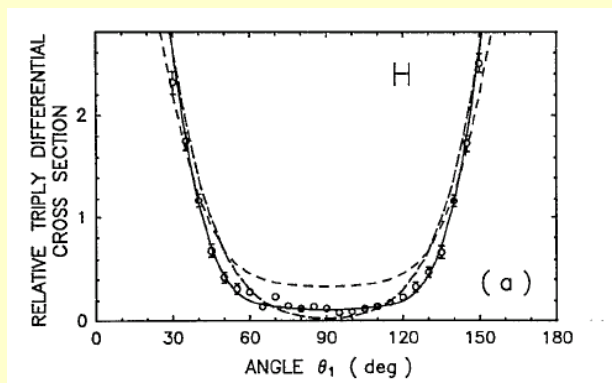
Kinematically Complete Electron – Impact Ionization

Ionization: Fragmentation of Atoms/Ions

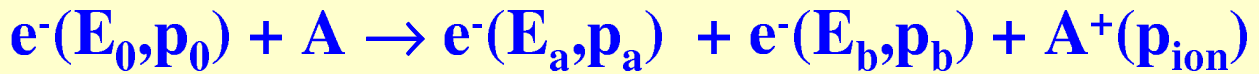


Comparison of 3-Body - Continua:

(Ehrhardt et al)



5-fold differential Ionization Cross Section:
5DCS for

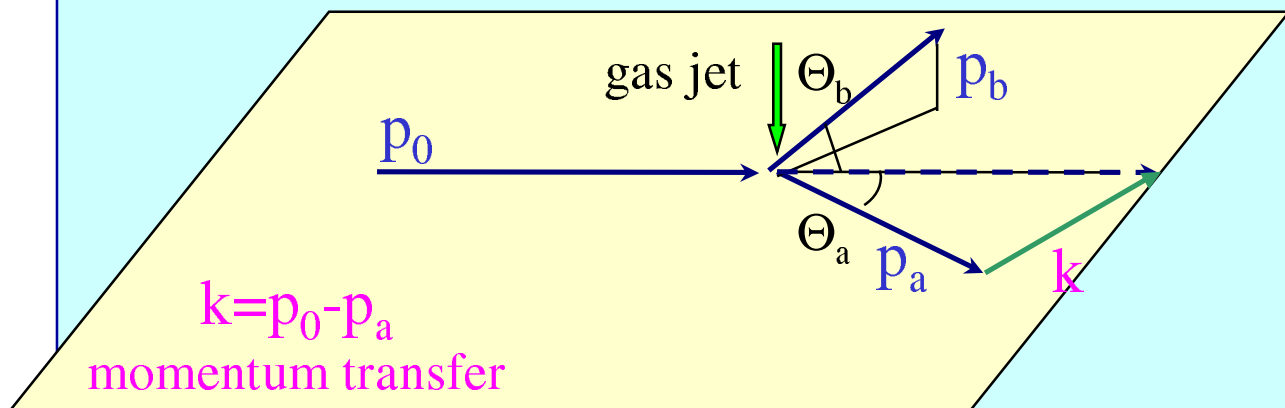


A kinematically complete **measurement** provides the 5-fold differential cross section

$$d^5\sigma(E_0, E_a, \Theta_a, \Theta_b, \varphi_{a,b})$$

5DCS : “COMPLETE” information;

➔ most stringent test of theories



peripheral collisions : IA with entire atom- : $k = \text{small}$
 binary “ : “ single e^- : $k = \text{large}$
 Quasi-Photoionization for $E_0 \gg E_{bind}$: $k \Rightarrow 0$

$$5DCS = \{\text{ionization-mechanism}\} \otimes \{\text{target-structure}\}$$

COMPLICATION:

♪ : in some areas of phase space these factors are coupled

What can be derived from Differential Cross Sections ?

DCS = {Ionization -mechanism} ⊗ {Target-structure}

Test of Theories:

DPW, CCC, ECS;
1. and 2. Order;

Wannier-threshold-
law;

WPR for Double-
Ionization;

relativistic DW;

Quasi-Photoionization

für $k \Rightarrow 0$;

1-Electron momentum-
density for individual
valence-orbitals;

many - body -
e-e effects in

$|i\rangle + |f\rangle$ states;

relativistic effects;

momentum dependence
of QED-Effects

in

vacuum-polarization,
selfenergy;

Why measure (e,2e) Ionization-cross sections for Ions?

1. study role of e-e interaction (e.g. correlation, polarization) via Z-dependence in DCS for H-, He-like systems
2. “The response of positive ions to ionizing radiation is a dominant process in the universe”
e⁻ -, Ion-impact- and Photo-ionization- cross sections for ions are needed recombination for radiative and rates in (see Opacity- project):
 - a) Fusionplasma
 - b) Astrophysical. Plasma, e.g. in Solar Corona and intra-galact. clusters

$\sigma_{\text{Ionization}}(\text{Electron - Ion})$

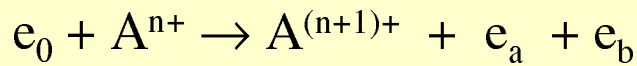
cannot be scaled from

$\sigma_{\text{Ionization}}(\text{Electron - Atom})$

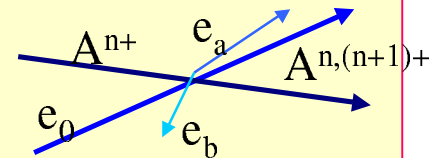
(Resonances / channel-couplings etc)

DCS and Inverse Kinematics

“crossed electron-ion beam” experiments

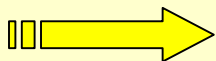
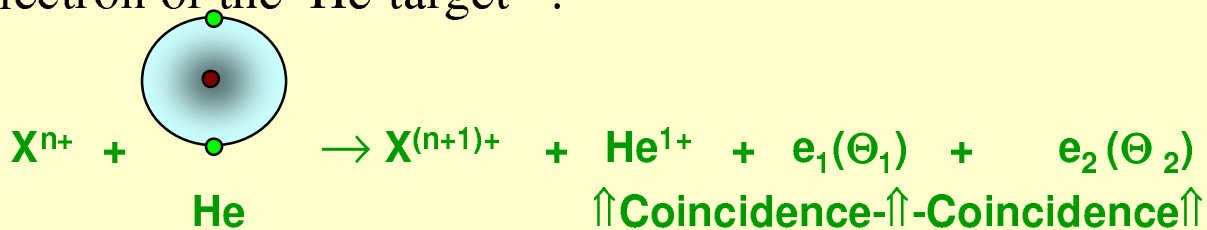


low target density $\rho \Rightarrow$ low luminosity L



$$L(e\text{-ion})/L(\text{ion-atom}) < 10^{-5}$$

Inverse kinematics: the incoming ion is ionized by an electron of the He-target :



REALM of REACTION-MICROSCOPE

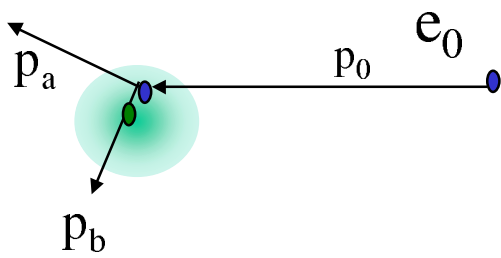
The He-target provides a dense source of quasi-free electrons

However: We must separate contributions:

ionization of X^{n+} by the He-nucleus (e-n)
from **ionization of X^{n+} via e-e interaction**

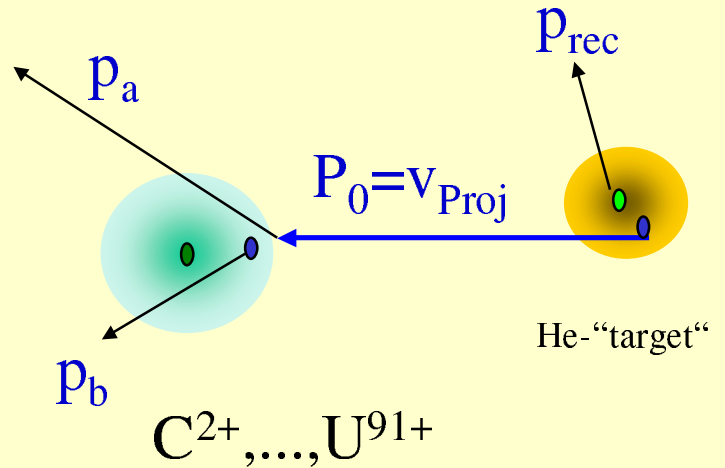
Electron Impact Ionization of Ions in inverse Kinematics

C^{2+}, \dots, U^{91+}



$E_0 = 0.5 - 200 \text{ keV}$

Projectile-system:

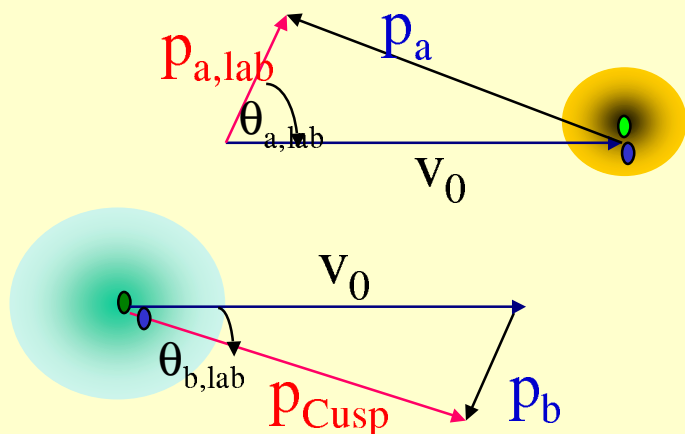


$p_B = \text{AI- electron or directly ionized}$

<u>Projectil-sys.</u>	<u>Labor.-sys.</u>
$p_{a,\text{gestr.}}$ fast	p_a slow
γ_A small	$\theta_{a,\text{lab}}$ large
$p_{b,\text{ion}}$ slow	$p_b \approx p_{\text{Cusp}}$
$0 < \gamma_b < 360$	$\theta_{b,\text{lab}}$ small

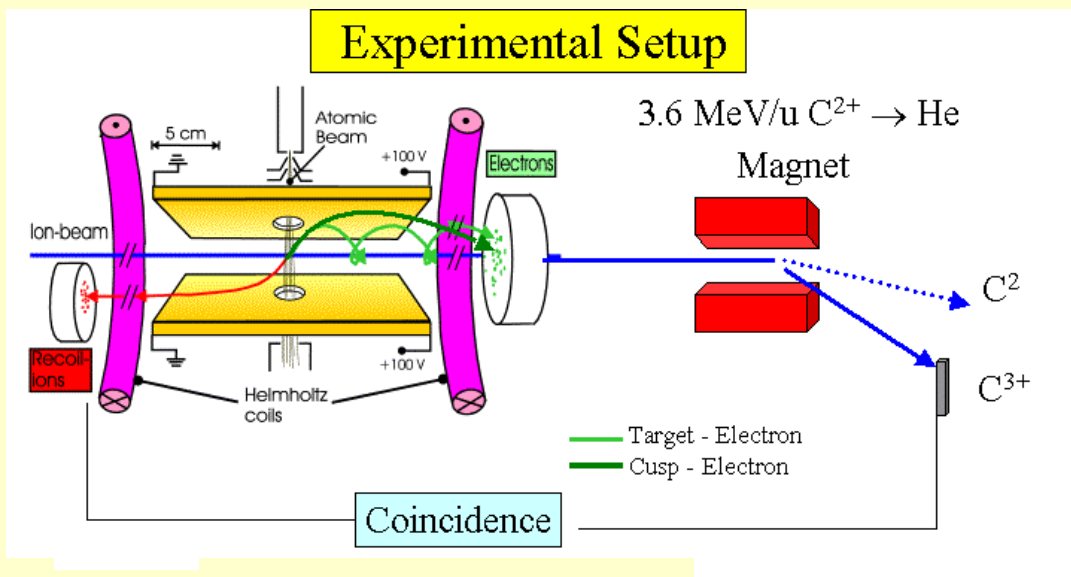
**\Rightarrow projectile electron
spectroscopy**

Laboratory-system:



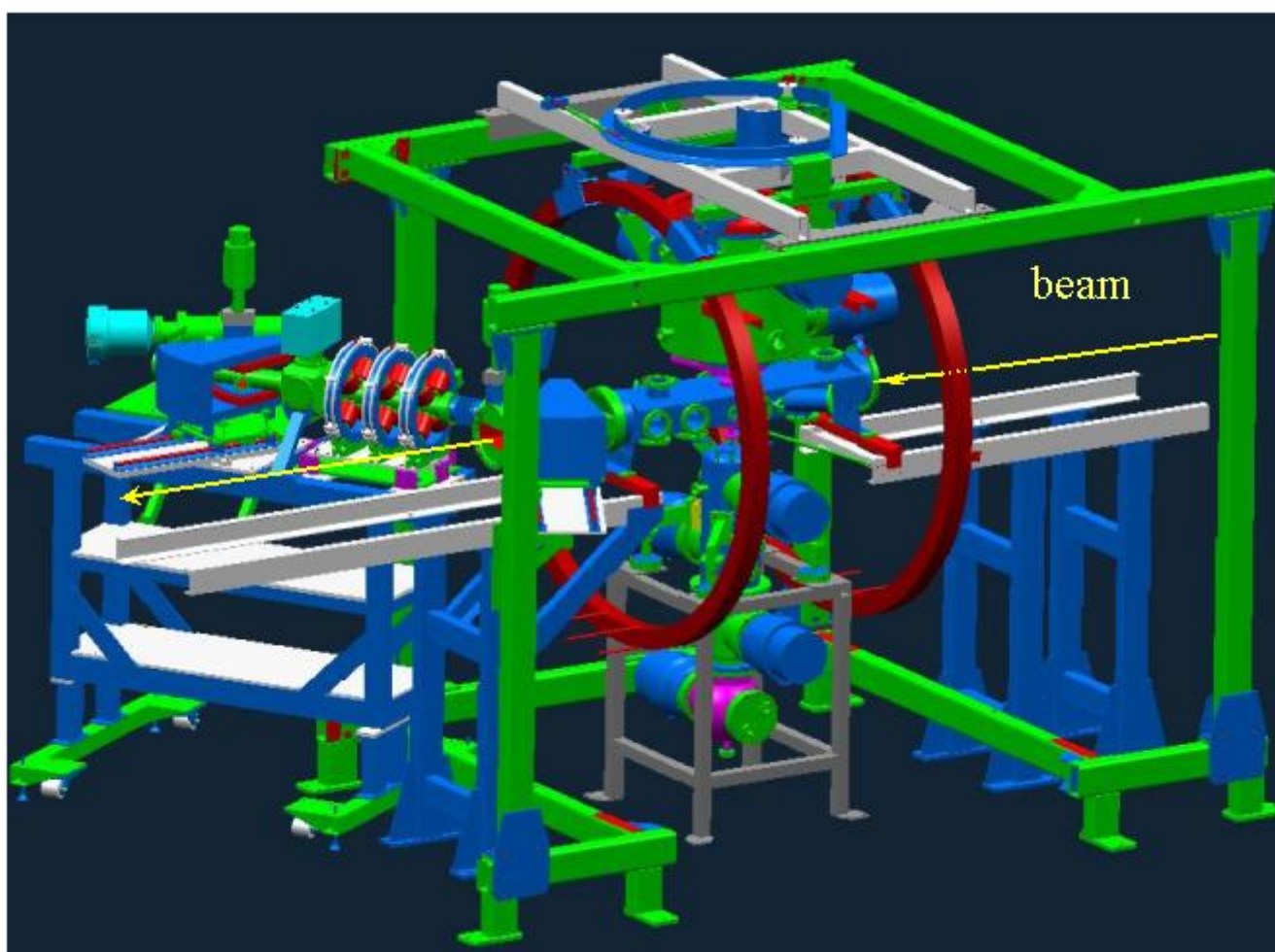
0.2 - 400 AMeV

Reaction microscope



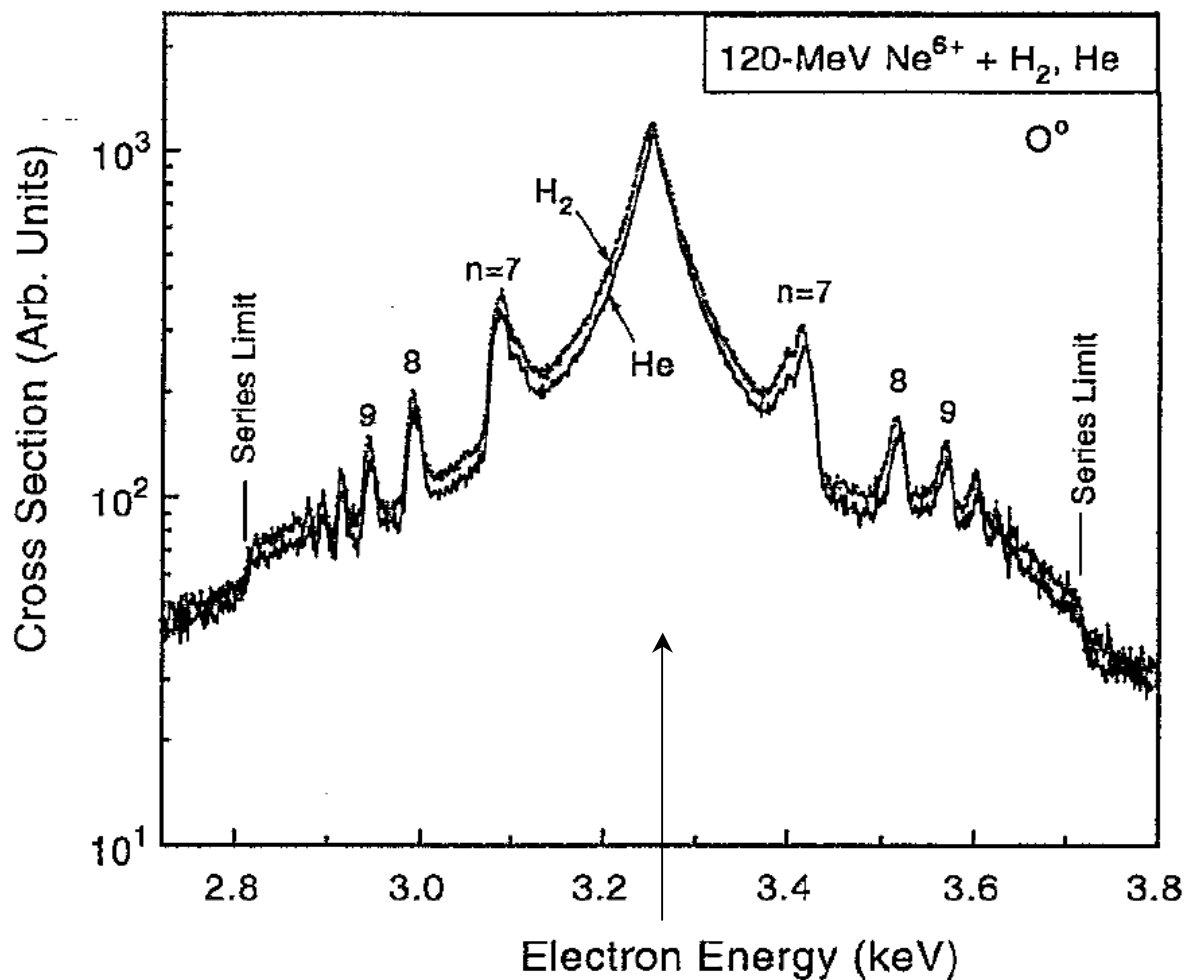
- For $v_{\text{Proj}} < 10$ a.u. both e^- can be mapped onto 2D position sensitive electron detector

Forward Electron Spectrometer and Reaction Microscope



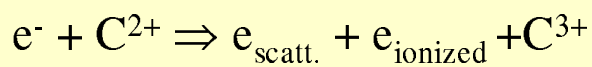
For high v_{Proj} the fast electron is analyzed with a separate spectrometer

Advantage of Fast Projectile Spectroscopy: Autoionization Spectra from Projectile Rydberg States



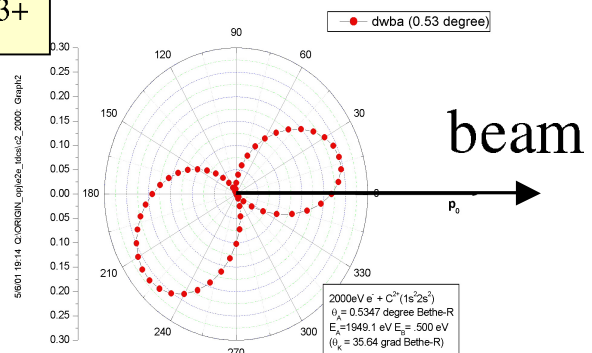
Stolterfoht et al

Electrons ionized into Projectile Continua

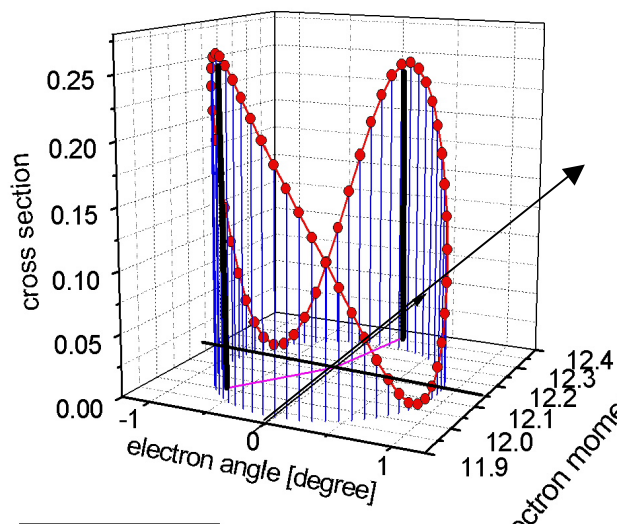


Projectile
frame

Theory: I. McCarthy



Laboratory
frame



Forward e⁻ Spectrometer

- 60° Dipol-QT-60° Dipol ,r=200mm, with aperture 100mm x250mm, for electrons with 20-200 keV, emitted into a narrow forward cone $2\Theta \leq 3^\circ$ around the beam axis:
- In telescope mode maps electrons from the target zone onto 2D-position sensitive detector with magnification $|M_x| = |M_y| = 1$; reconstruction of scattering plane and primordial direction of emission following momentum analysis, variable momentum resolution modes
- In high resolution mode $p/\Delta p = 1600$

*First Experiments at JRM :
5DCS at threshold in
 $\Theta_{a,b} = \text{constant}$ geometry*

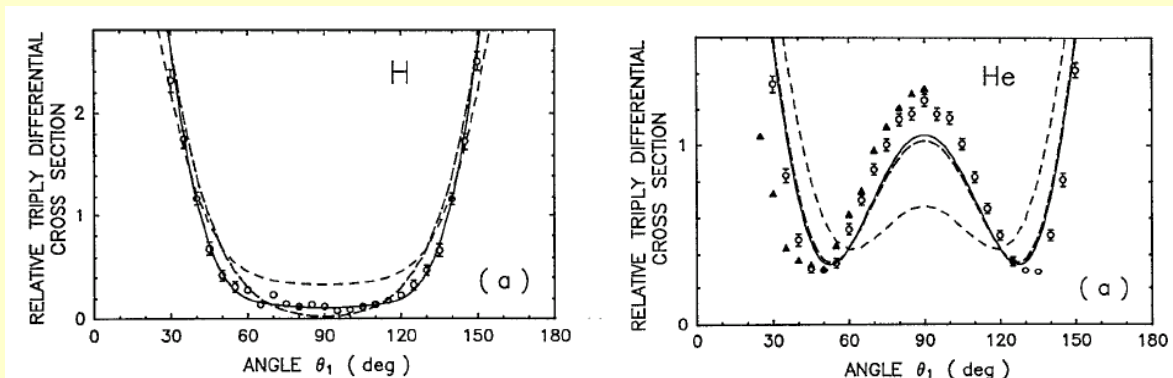
for $E_0 \geq \epsilon_{\text{Bind}}$, close to threshold:

3-body- effects in entrance channel:
(polarization of the targets via incoming
electron) and PCI cannot be distinguished
in asymmetric geometry;

for $\Theta_{ab} = \text{const.}$ and $E_a = E_b$ the effect from
the entrance channel can be isolated:

e.g. : $E_a = E_b = 2\text{eV}$

Ehrhardt et al.



Quasi-Photo-ionization of FeXVII: Opacity Project for Stellar Atmospheres

- Goal: determine the Equation of State for stellar atmospheres for $T \leq 10^7$ K
- EOS must reliably predict the optical properties of the plasma (current models derive from the observed spectra widely varying T-distributions)
- To improve models:
- Reliable input-data are missing :
 - a) Ionization- Rekombination over a wide q-range
 - b) **Photoionization of ions**
 - c) Dielectronic Recombination

*(e,2e) at the ESR:
Photoionization of Fe XVII via Electron-
Electron-Coincidences*

Motivation: Opacity-Project

Knowledge of Fe ion distribution is of fundamental importance in modeling stellar atmospheres (and intra-Cluster clouds)

Fe is dominant source of opacity over a wide range of high temperatures in stellar atmospheres.

Photoionization contributes $\approx 50\%$ or more for opacity of the Fe .

Hibbert et al. pointed to large discrepancies between non-relativistic and relativistic

Photoionization cross sections for Fe XVII
(PRA57,3489(1998))

below the $2s^22p^5$ ($^2P_{1/2}$)
threshold ,where a large
fraction of the oscillator-
strength is.

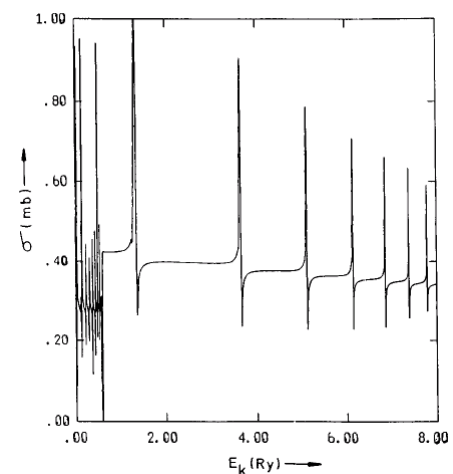


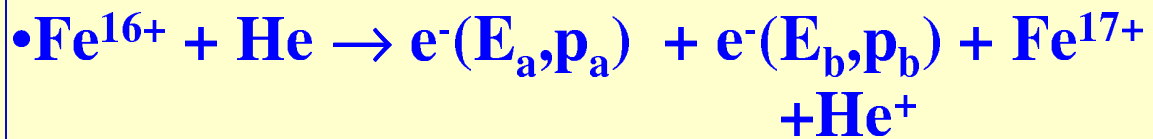
Photo-Ionization via Electron-Electron Coincidence

- For swift, peripheral collisions and small momentum transfer ΔK :
- $DCS(e^- + Atom) = DCS(h\nu + Atom)$
the interaction of a swift electron with an atom is equivalent to the interaction with an incident photon, the target electron experiencing a force along the polarization vector of the light ; [provided that $E_{inc} \gg E_{bind}$]:

$$E_{loss} = E_0 - E_a = E_{h\nu}$$

- For small ΔK the generalized oscillator – strength for ionization is given by the optical oscillator strength (Bethe)

Photo-ionization of Fe XVII via Electron-Electron Coincidence



Coincidences between outgoing products

$e_{\text{fast}}, e_{\text{slow}}, \text{He}^+$ und Fe^{17+} for $E_{\text{proj}} \geq 100 \text{ AMeV}$.

• Using the new combined reaction-microscope \oplus magnetic 0° -degree electron-spectrometer

$$E_{0,\text{inc}} - E_{\text{scatt}} = E_{\text{hv}}$$

→ **Energy of the equivalent
incident photon;**

$E_b, \Theta_b [e_{\text{fast}}]$ around 0° Cusp -->
complete angular distribution for photo-
electrons on 2D position sensitive
electron-detector