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**Trapping and Manipulating Atomic and Subatomic Particles** 

# **Principles of (Ion) Storage Rings**

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## Outline

- 1. Why do we need ion storage-cooler rings or ion traps? Traps and Rings - complementary devices
- 2. Basic concepts of rings: ion optics, storage equation of motion
- 3. Beam-cooling techniques
- 4. Tools for the diagnosis of stored and cooled beams

#### 1. Why do we need storage-cooler rings or traps?



#### Trifodi nebula, glowing in the infrared light of a nascent star

# The major part of matter in the universe exists as a plasma, in a well-balanced equilibrium of ions and electrons

Ion storage rings (traps): doing atomic physics, nuclear physics and astrophysics with ions in wellpreserved atomic charge states *Highly charged ions* can be *stored* for hours and *'cooled'* in *Storage-Cooler Rings* where their properties are investigated by *sophisticated methods* 



*dense plasma of ions and e<sup>-</sup>* (10<sup>26</sup>/cm<sup>3</sup>)

*high temperature* kT ≈ 1...10 keV

*excited* atomic and nuclear *states* 

*equilibrium* of charge states

highly charged ions low density (10<sup>4</sup>/cm<sup>3</sup>)

**low temperature kT \approx 0.1 eV** (T= 10<sup>3</sup>K)

ions in ground-state

**non-equilibrium** of charge states

**Ion Traps** 

localized

fixed (low) velocity

medium charge states

few...one ion

**Ion Storage Rings** 

extended

tunable high velocity

high charge states

many...one ion

cooling techniques

restricted lifetime

photonic excitation

radioactive decay

mass determination

lifetime measurement

## 2. Basic concepts of storage rings: ionoptics, storage, equation of motion

To store ions one needs a set of bending and focusing magnetic multipoles, a 'lattice'.



**Experimental Storage Ring ESR at GSI** 

Due to the finite emittance ε ( = size x divergence) ions move on *quasi-periodic orbits* performing horizontal and vertical 'betatron oscillations' around the 'reference orbit'

with a 'betatron wavelength'  $[\lambda_{\beta}]_{h,v}$ 



[orbit length C] /  $[\lambda_{\beta}]_{h,v}$  = 'tune  $Q_{h,v}$ '

The orbits in terms of deviations from the 'Sollbahn' (including 'orbit dispersion' due to  $\Delta p \neq 0$ ) are solutions of 'Hill's differential equation', first proposed by G.W.Hill in 1886 to describe the motion of planets disturbed by 3-body forces (for a Paul trap it's the Mathieu equation):

**u''** (s) + k(s)u -  $1/\rho \Delta p/p = 0$ 

u(s) = deviation from reference orbit; s = ion path; k(s) = k(s + C)
periodic force, determined by the 'lattice'; C = circumference

*'cos-like' solutions* without 'dispersion' ( $\Delta p = 0$ ):

$$\mathbf{u}(\mathbf{s})_{\mathbf{h},\mathbf{v}} = \{ \varepsilon_{\mathbf{h},\mathbf{v}} \ \boldsymbol{\beta}_{\mathbf{h},\mathbf{v}}(\mathbf{s}) \}^{1/2} \ \cos\{ \ \boldsymbol{\psi}(\mathbf{s}) + \delta \ \}$$
$$\mathbf{'phase'} \ [\boldsymbol{\psi}(\mathbf{s})]_{\mathbf{h},\mathbf{v}} = \int_{0}^{s} d\mathbf{s'} / \boldsymbol{\beta}(\mathbf{s'})_{\mathbf{h},\mathbf{v}}$$

then the 'betatron amplitudes'  $\beta(s)$  depend only on k(s) 'Strong focusing' gets small  $\beta$ , large  $\psi$ 

with 'betatron wavelength'  $[\lambda_{\beta}]_{h,v} \equiv 2\pi C/\psi(s = C)_{h,v}$ 

$$\rightarrow$$
 'tune'  $Q_{h,v} \equiv C / [\lambda_{\beta}]_{h,v} = \psi(s = C)_{h,v} / 2\pi$ 

Avoid integer (n) or 'algebraic' (m/n) Q<sub>h,v</sub>

It generates periodic, resonance-like enhancement of small perturbations and immediate loss of beam



Such a 'Q- resonance' -due mainly to Jupiter's periodic force- forced the destruction of a former planet orbiting between Jupiter and Mars. Today still 'gaps' (e.g. Hecuba) exist.

Ion storage rings have a restricted acceptance given by  $\varepsilon_h$ ,  $\varepsilon_v$ , (typically 10...30  $\pi$  mm mrad) or  $\Delta p/p \approx 1...5$  %.

These values and the set of magnets (the 'lattice') determine the solution of Hill's equation, i.e. the actual betatron wavelengths, the betatron amplitudes  $\beta(s)_{h,v}$  and  $Q_{h,v}$ .

**Q** depends on momentum spread  $\Delta \mathbf{p}$  and beam current:

 $\Delta Q$  large  $\rightarrow$  crossing resonances  $\rightarrow$  beam loss



# **Q changes** too, when **accelerating/decelerating** ions

 $\rightarrow$  fast crossing of resonances needed

# **Coulomb repulsion** of ions with charge q effects a **Space-charge limit** of number N of ions ( $\Delta p = 0$ ):

$$N_{max} < \pi/r A/q^2 β^2 γ^3 ε_{-} [1 + {ε_{+} Q_{-} / ε_{-} Q_{+}}^{1/2}] ΔQ$$

{ r = classical proton radius,  $1.5 * 10^{-18}$  m;  $\varepsilon_{+/-}$ ,  $Q_{+/-}$  = larger/smaller of emittances (tunes);  $\Delta Q$  = 'tune shift' ( < 0.05...0.2)}

**Examples:** (emittances  $\varepsilon$  in [ $\pi$  mm x mrad])

1. 
$$\mathbf{U}^{92+}$$
 (Q<sub>+</sub> = Q<sub>-</sub> = 2,43;  $\Delta \mathbf{Q} = 0.05$ ) :

- **not** cooled,  $\varepsilon_{+/-} = 1$  **cooled**,  $\varepsilon_{+/-} = 0.1$
- $E = 500 \text{ A MeV} \rightarrow N_{\text{max}} \approx 4 \text{ x } 10^{12} \qquad N_{\text{max}} \approx 4 \text{ x } 10^{11}$

 $E = 50 \text{ A MeV} \rightarrow N_{max} \approx 2.2 \text{ x } 10^{11} \text{ N}_{max} \approx 2.2 \text{ x } 10^{10}$ 

### 2. Ne<sup>10+</sup>

$$E = 500 \text{ A MeV} \rightarrow N_{max} \approx 3.4 \text{ x } 10^{14} \qquad N_{max} \approx 3.4 \text{ x } 10^{13}$$
$$E = 50 \text{ A MeV} \rightarrow N_{max} \approx 1.9 \text{ x } 10^{13} \qquad N_{max} \approx 1.9 \text{ x } 10^{12}$$

in addition: **'chromaticity'** (momentum dependence)..... 11

# 3. Beam-cooling techniques

# 'Cooling' :

# **Enhancing the phase-space density**

# i.e. Reducing size, $\Delta p$ , $\rightarrow$ emittance $\varepsilon$

### but:

for **fixed velocity v**, ε is constant (Liouville)



## We can never change the area = $\pi \epsilon$ of the ellipse

 $\rightarrow$  apply external interactions  $\rightarrow$  cooling

# for ion storage rings three cooling methods: 'stochastic'-, 'electron'- and 'laser'-cooling



Simon van der Meer 1925\*, CERN, Nobel award 1984 invented 1983 'stochastic cooling' to get a cooled antiproton beam in the CERN SPS collider for detection of the W Boson the 'mediator' of weak interaction



#### Principle of 'stochastic' cooling [www.fz-juelich.de/ikp/ikp-general/images/stochast]

pick-up:

→ beam position, deviations from reference orbit

**difference** signal (left/right)  $\rightarrow$  amplifier

 $\rightarrow$  kicker at (n + 1/4)  $\lambda_{\beta}$ 

cooling time  $\tau_{sto} \propto N_{ion}$  /Bandwith;  $\Delta p/p \approx 10^{-3}$ 

#### laser cooling proposed first by Ted Haensch realized in ion storage rings first at TSR and ASTRID



Test Storage Ring TSR Heidelberg

- 1. **resonant absorption** of a photon  $E_{\gamma}$
- 2. directed momentum transfer on ion,  $\Delta p = E_{\gamma}/c$
- 3. isotropic re-emission of  $E_{\gamma} \rightarrow$  net momentum transfer
  - $\rightarrow$  shifting the velocity profile  $\Delta v$  of the stored ions

 $\rightarrow$  second counteracting device to narrow  $\Delta v$ 

→ **interaction** of **longitudinal and transverse** d.o.f. 15



The counteracting force to narrow ∆p may be a second laser or -as at the TSR- a high voltage tube

laser cooling realized for Be<sup>+</sup>, Li<sup>+</sup>....

for highly charged ions (H-like, He-like) only chance to use the 1s hyperfine splitting

advantage of storage rings:

resonance Lorentz-boosted

 $\lambda_{rf} = \lambda_0 \{ (1 - \beta)/(1 + \beta) \}^{1/2} (laser 180^{\circ})$ 

#### electron cooling invented in the sixties by G. Budker



#### collinear, 'cold' (sharp velocity) electrons, guided by a magnetic field, exchange longitudinal and transverse momenta with the ions by Coulomb collisions

electrons removed after one turn

ions come back after  $\approx 1 \mu s$ 

# ions get the electron temperatures T<sub>1</sub>, T<sub>t</sub> after equilibrium has been reached:

$$\mathbf{kT}_{1} = \mathbf{A} \ \beta^{2} \ \mathbf{c}^{2} \ (\Delta \mathbf{p}/\mathbf{p})_{e}^{2}$$
$$\mathbf{kT}_{t} = 1/2 \ \mathbf{A} \ \gamma^{2} \ \beta^{2} \ \mathbf{c}^{2} \ \{(\epsilon/\beta)_{h} \ + \ (\epsilon/\beta)_{v} \ \}_{e}$$

Time  $\tau_{\rm C}$  needed for cooling:

 $\tau_C \propto A/Z^2 \ 1/j_{e-} \ \Delta v \ ^3$ 

[0.1s...30s; A = 200, Z=80; je- = 200 mA)

 $[\Delta v/v]_{inj} = 10^{-3}....10^{-1}$  (hot fragments);  $j_{e} = 200 \text{ mA}$ 

#### **Electron Cooling** provides:

• brillant beams ( $\varepsilon = 0.1 \pi$  mm mrad;  $\Delta p/p = 10^{-6}$ )

• constant energy

- long storage times (hours for few-e<sup>-</sup> ions)
  - same velocity for all ions



Theoretical cooling force:  $\propto 1/\Delta v^3$ 





Horizontal beam size before and after electron cooling (TSR)

## Since cooled ions all have (almost) the same velocity

a one-to-one correspondence between circulation frequency f and m/q- ratio exists

### 4. Diagnostic tools: 'Schottky noise'





 $\rightarrow$  df/f = d $\beta/\beta$  - d $\rho/\rho$ 

 $\rightarrow$  df/f = { 1/ $\gamma^2$  -  $\alpha_p$  } dp/p

df/dp = 0 for  $\gamma_t^2 = 1/\alpha_p$ ,  $\gamma_t$  = 'transition energy'

 $\alpha_p$  = 'momentum compaction'(property of lattice) =  $(d\rho/\rho)/(dp/p) \approx 1/7$  (ESR)

for v, B = const.  $\rightarrow \rho \propto m/q$ {f<sub>1</sub> - f<sub>2</sub>}/f<sub>1</sub> = - $\alpha_p$  {[(m/q)<sub>1</sub> - (m/q)<sub>2</sub>] / [(m/q)<sub>1</sub>]}



#### Test Storage Ring TSR at the MPI Heidelberg www.mpi-hd.mpg.de/be/tsr/tsrindex



Experimental Storage Ring ESR at GSI Darmstadt www.gsi.de/accelerator/esr



#### **Electron cooler at the ESR** www.gsi.de/accelerator/esr