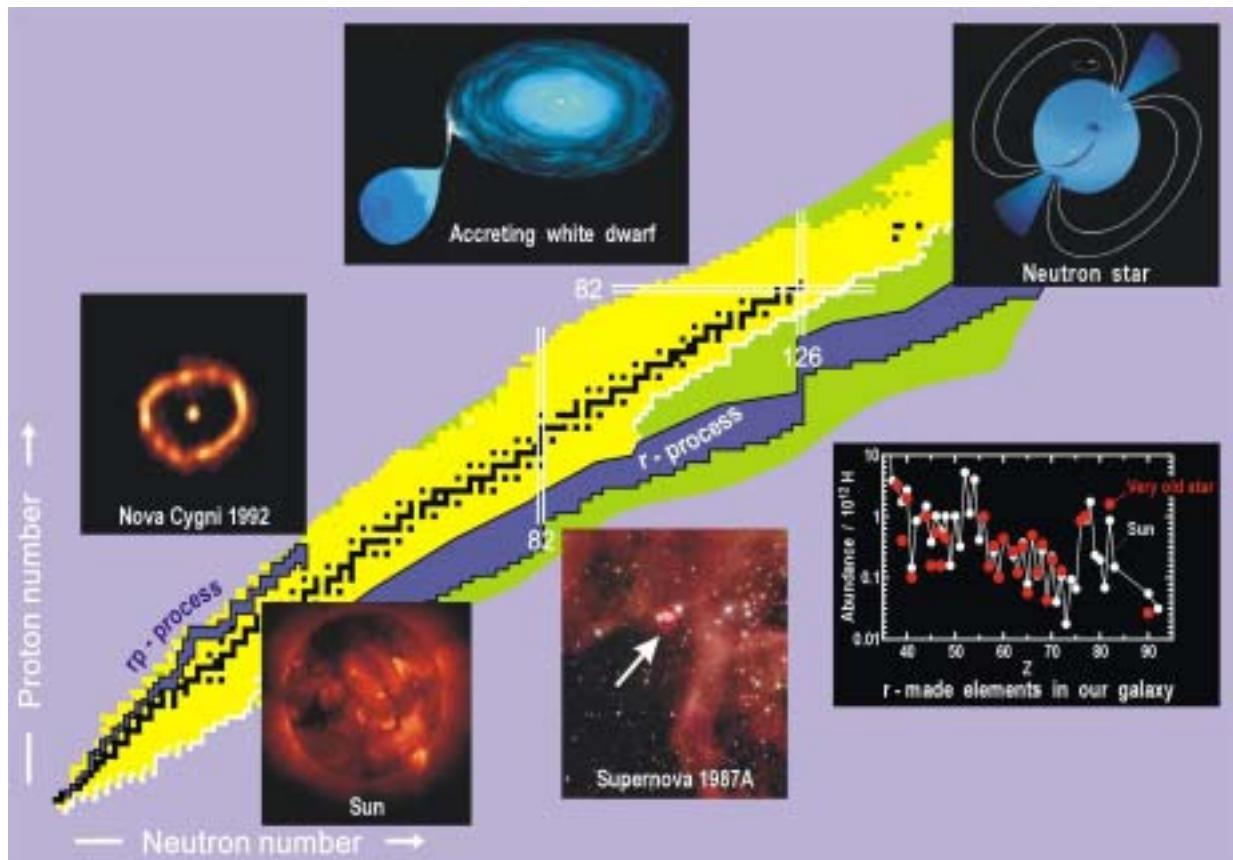


# Impact of atomic charge state on stellar nucleosynthesis

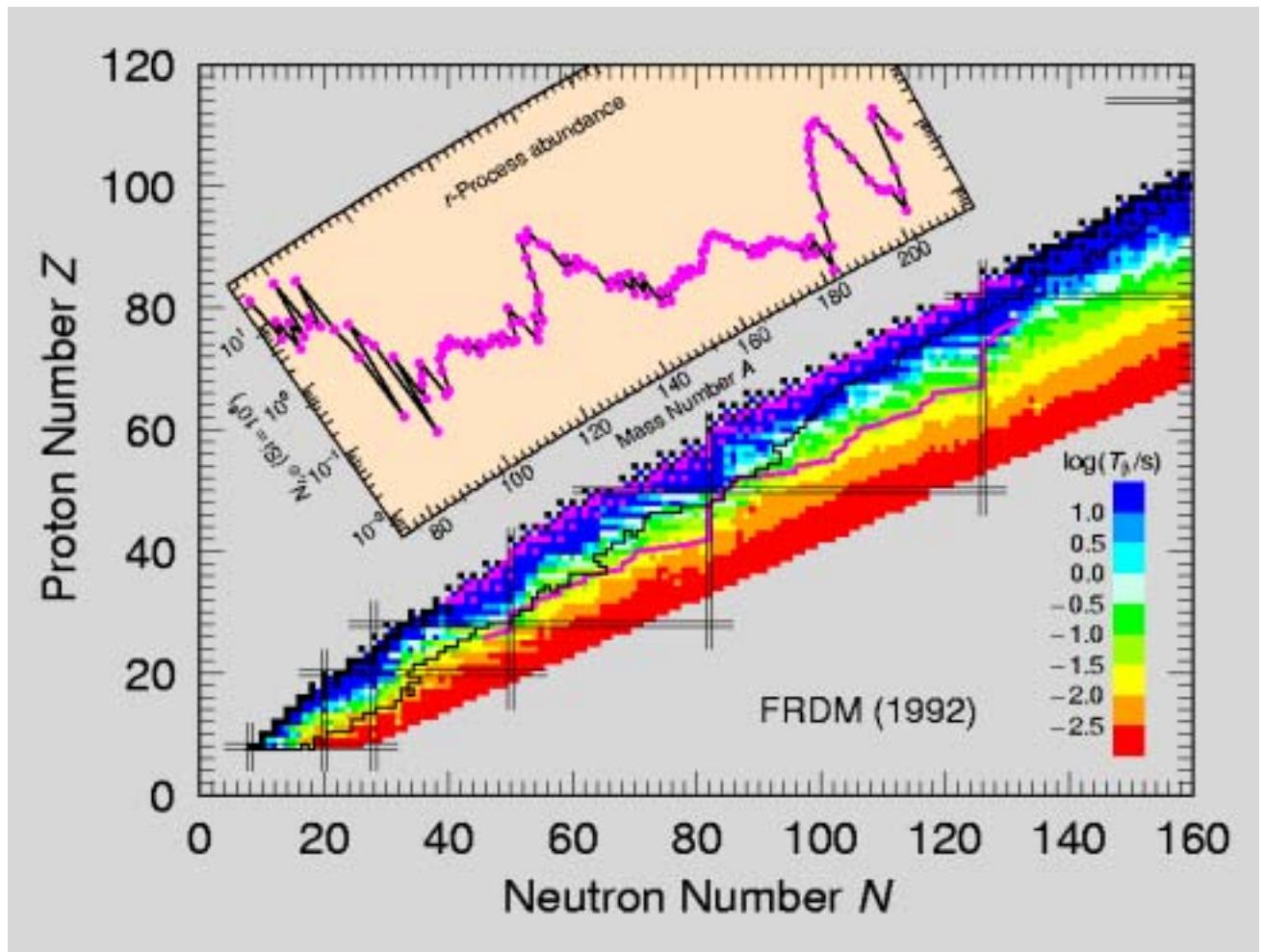
Fritz Bosch, GSI



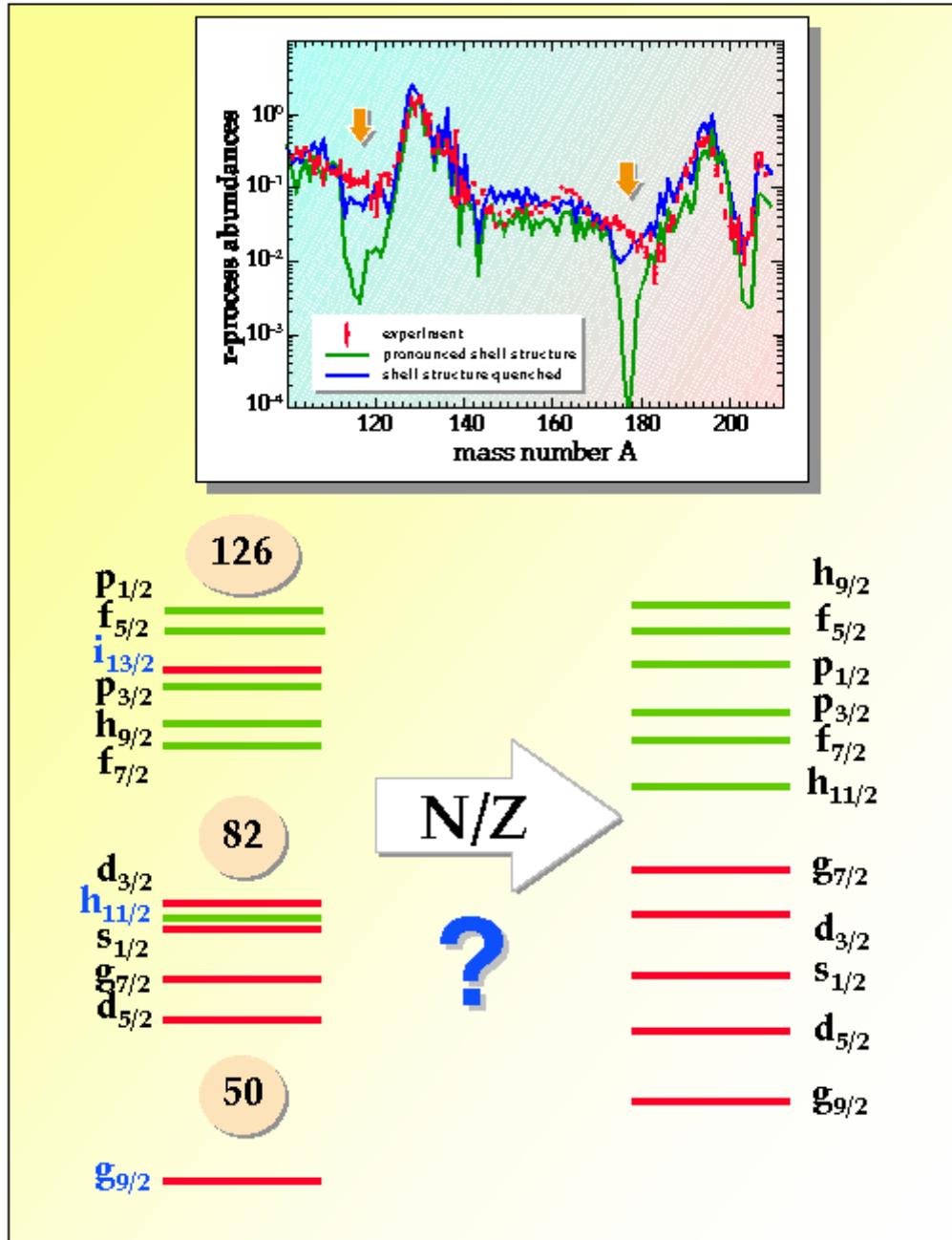
1. Stellar nucleosynthesis (besides fusion) driven by neutron (proton) capture,  $\beta^-$  and  $\beta^+$  decay at high temperatures = high atomic charge states
2. masses and  $\beta$  lifetimes determine pathway of nucleosynthesis and abundance of elements
3.  $\beta^-$ ,  $\beta^+$  lifetimes depend on atomic charge state
4. storage rings only tools to measure  $\beta$  decay of HCl

# Explosive nucleosynthesis ( $r$ , $rp$ ) *not yet understood*

$r$ -process(es): rapid neutron capture  
one or several, site, pathway?



→ abundance of heavy elements in our solar system

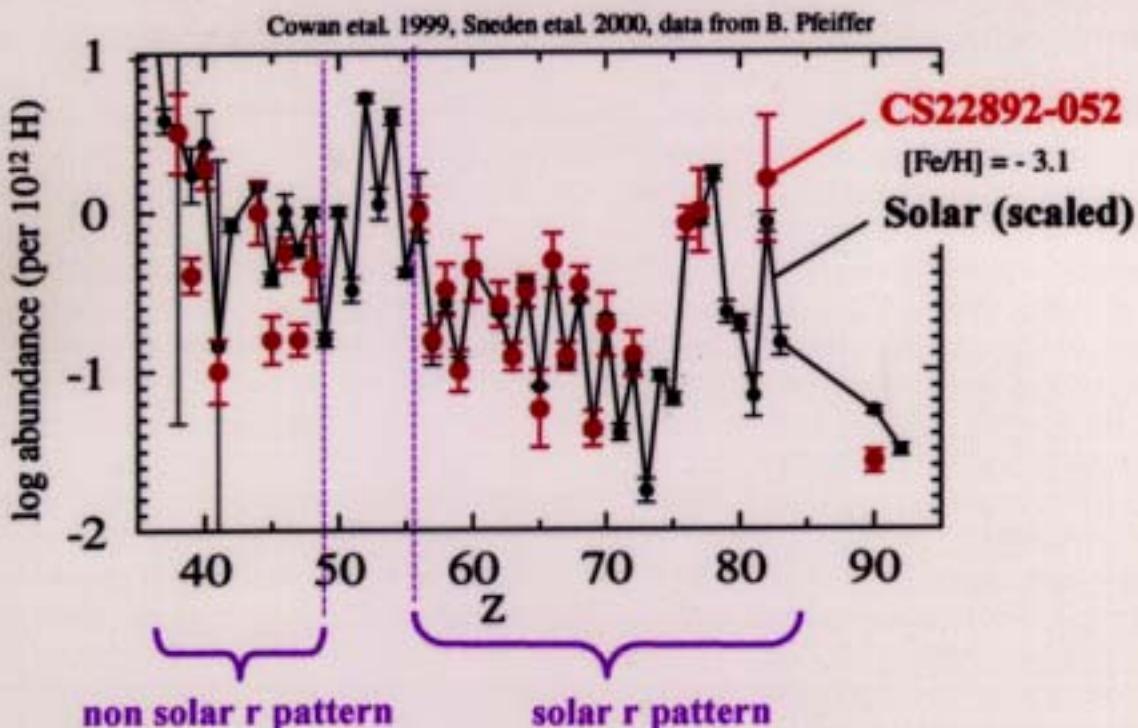


## Entanglement of nuclear structure and nucleosynthesis in the r-process Quenching of shell gaps?

→ produce **neutron-rich, highly charged atoms**  
and measure their ground state properties:  
**masses and  $\beta^-$  lifetimes**

## Ultra metal poor halo stars

A window into the composition of the galaxy at age 0.1-0.7 Gyrs



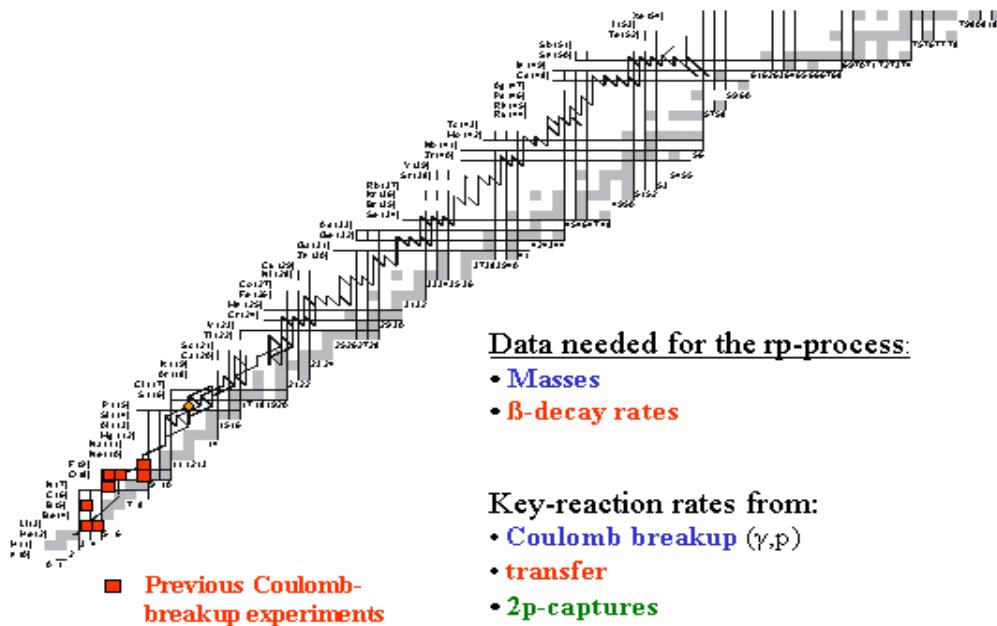
### Open questions:



What are these r process events ?

Two r processes ?

## *The rp-process*



### Data needed for the rp-process:

- Masses
- $\beta$ -decay rates

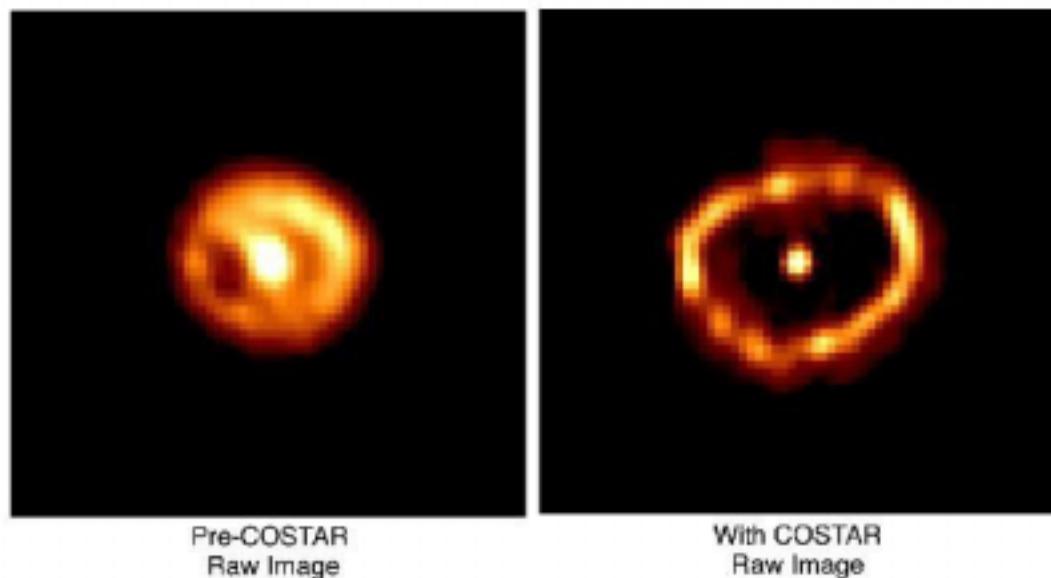
### Key-reaction rates from:

- Coulomb breakup ( $\gamma, p$ )
- transfer
- 2p-captures

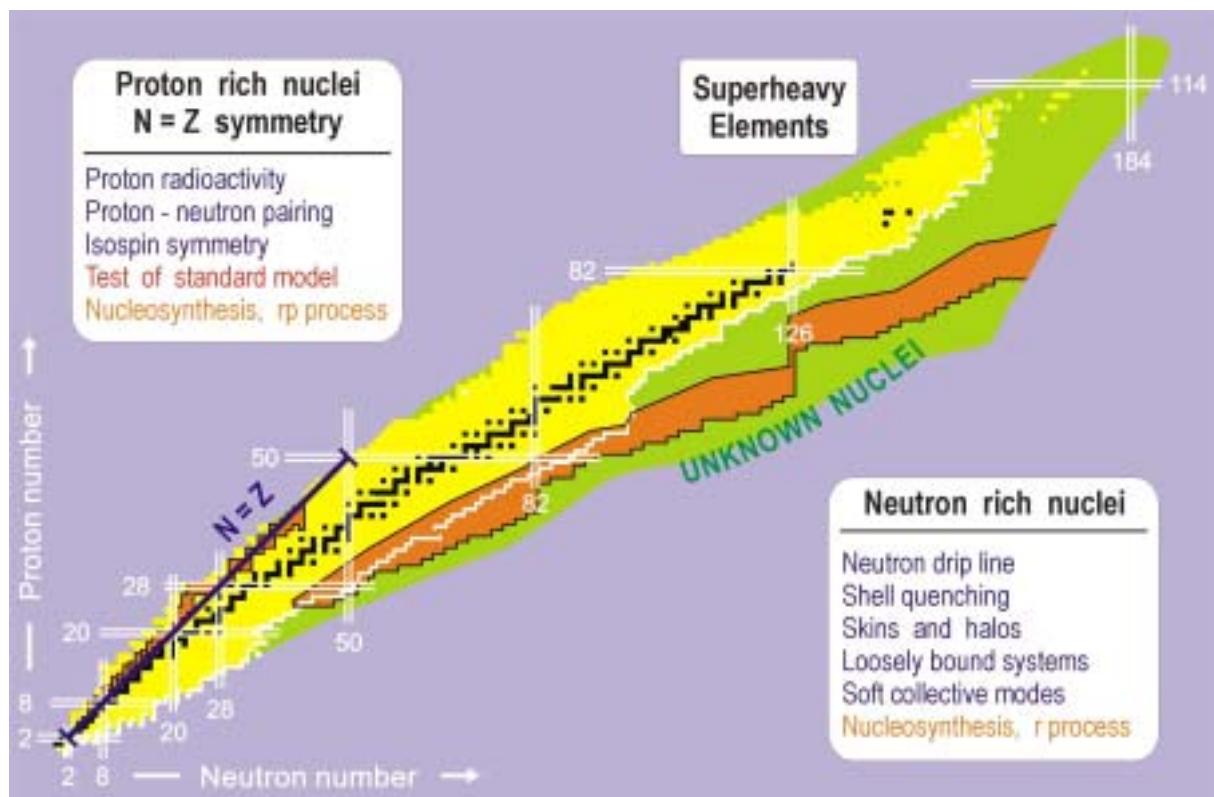
**rp (rapid proton capture) -process:  
explosive H-burning on the surface of mass  
accreting white dwarfs at high T**

**ground state properties of p-rich, highly-charged atoms:  
masses and  $\beta^+$  lifetimes**

**Nova Cygni 1992**  
Hubble Space Telescope  
Faint Object Camera



**Nova outburst signals explosive H burning up to tin along the proton dripline via p-capture, p-decay,  $\beta^+$  decay**



# Why and how $\beta$ lifetimes depend on atomic charge state

a) orbital electron capture EC:  $\lambda_{EC} = f(n_{eb})$

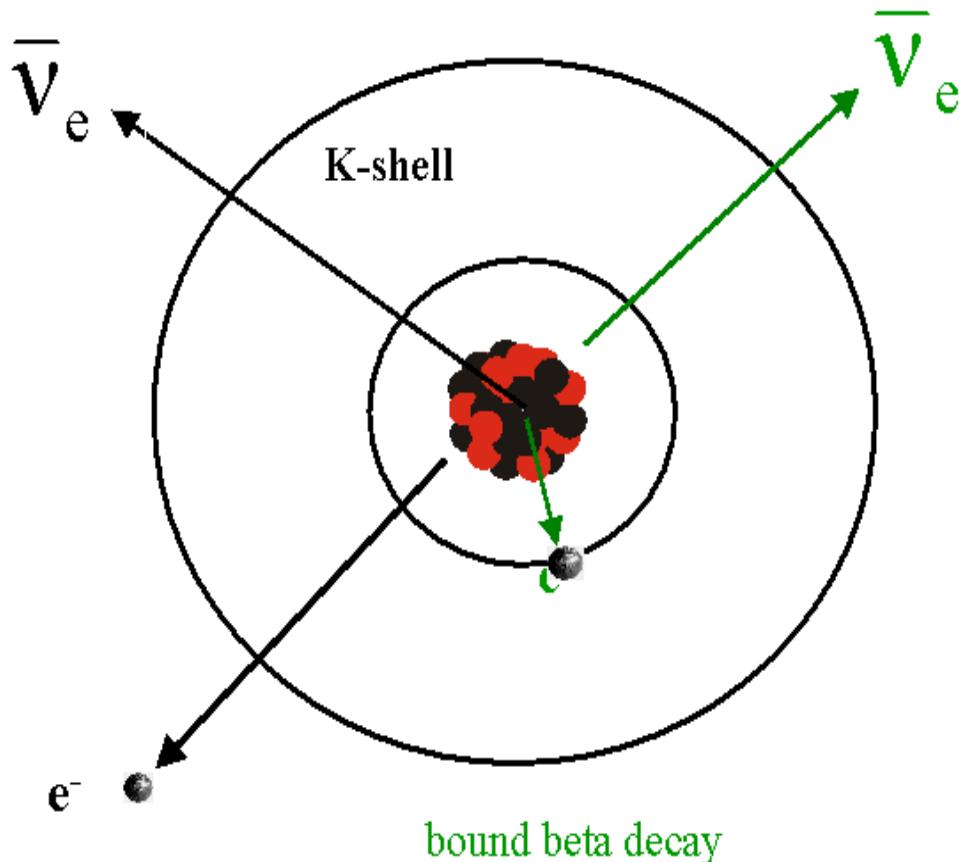


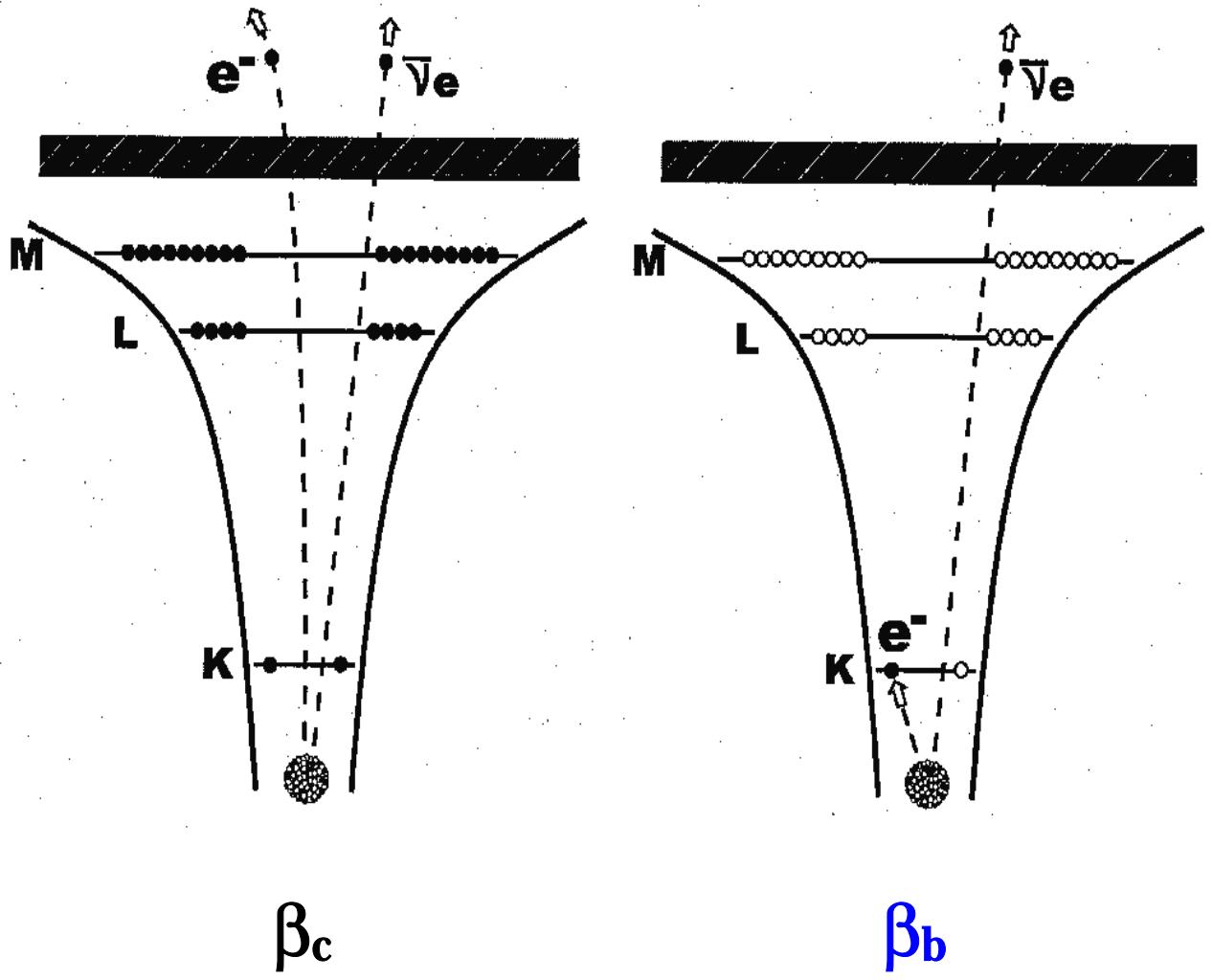
if  $Q_{EC} < 2 m_e c^2$  atom gets **stable** for  $n_{eb} = 0$

b) bound beta decay  $\beta_b$  = time-mirrored EC



if inner shell **vacancies** available





**Q value larger**  
**by  $e^-$  binding energy**

decay probability  $\lambda_{\beta b} > \lambda_{\beta c}$

stable neutral atoms get **unstable**

if **Q(neutral) < |E<sub>b</sub>(K)|**

The couple  $^{163}\text{Dy}/^{163}\text{Ho}$  ( $^{205}\text{Tl}/^{205}\text{Pb}$ ,  $^{193}\text{Ir}/^{193}\text{Pt},\dots$ )

a) *neutral* atoms (on earth)

$$\begin{array}{c} ^{163}\text{Ho}^{0+} \\ \hline \\ \text{T}_{1/2} = 4570 \text{ y} \end{array}$$

$$\begin{array}{c} ^{163}\text{Dy}^{0+} \\ \hline \\ \text{T}_{1/2} = \infty \end{array}$$

b) *bare* nuclei in stellar plasma

Cont.

$$\begin{array}{c} ^{163}\text{Dy}^{66+} \\ \hline \\ \text{T}_{1/2} = 47 \text{ d} \end{array} \quad \begin{array}{c} \hline \\ \text{----- L} \end{array}$$

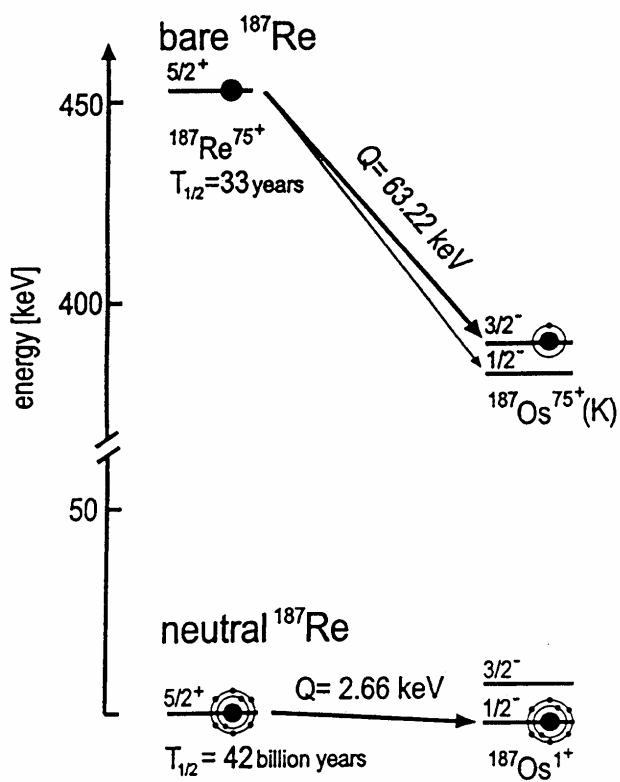
$$Q_{\beta b}^K = + 50 \text{ keV}$$

$$\begin{array}{c} ^{163}\text{Ho}^{66+} \\ \hline \\ \text{T}_{1/2} = \infty \end{array} \quad \begin{array}{c} \hline \\ \text{----- K} \end{array}$$

$$\begin{array}{ccccccc} Q_{\beta b}^{K,L} = & - Q_{EC} & - & \Delta B_e & + & B_{K,L} & \\ & - 2.6 & & - 12.5 & & + 65.1 & = + 50 \text{ keV (K)} \\ & & & & & + 16.3 & \approx + 1 \text{ keV (L)} \end{array}$$

# The couple $^{187}\text{Re}/^{187}\text{Os}$ ( $^{205}\text{Tl}/^{205}\text{Pb},\dots$ )

$\beta_b$  decay goes to an **excited state, forbidden in  $\beta_c$**



$\beta_b$  decay provides  
first possibility to measure the  
branching ratio  $\lambda_{\beta b} / \lambda_{\beta c}$   
for the *same* nuclear transition

$$\lambda_{\beta b} / \lambda_{\beta c} = [Q_{\beta b} / Q_{\beta c}]^2 \psi(0)^2 / f(Z, W_0)$$

$\Rightarrow$  *Fermi function*  $f(Z, W_0)$   
*for  $\beta^-$  decay*

still missing complement to the well known  
*Fermi function for  $\beta^+$  decay* from  
 $\lambda_{\beta+} / \lambda_{EC}$

### c) internal conversion IC (isomers)

for high multipolarities (E4, M4...) IC large

lifetime much longer (10,...10<sup>6</sup>) without e<sub>b</sub>

neutral <sup>193</sup>Ir<sup>0+</sup>

T<sub>1/2</sub> = 10 d \_\_\_\_\_ <sup>193</sup>Ir\*<sup>0+</sup> 11/2<sup>-</sup> E = 80 keV

M4

\_\_\_\_\_ <sup>193</sup>Ir<sup>0+</sup> 3/2<sup>+</sup>

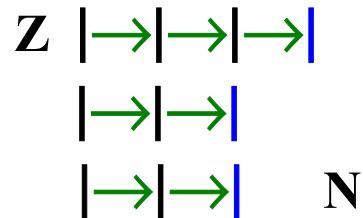
bare <sup>193</sup>Ir<sup>77+</sup>

T<sub>1/2</sub> > 10<sup>3</sup> y \_\_\_\_\_ <sup>193</sup>Ir\* 77+

M4 (no e<sub>b</sub>)

\_\_\_\_\_ <sup>193</sup>Ir<sup>77+</sup>

## s (slow neutron capture)- process:



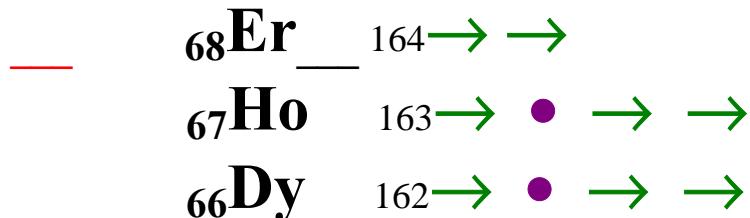
near the valley of stability: n-capt. / $\beta^-$  decay  
 $\beta^-$  decay for  $T_{1/2}(\beta) \leq 1$  year

neutron density  $\approx 10^8 \text{ cm}^{-3}$   
 electron density  $\approx 10^{27} \text{ cm}^{-3}$   
 temperature  $T_S \approx 3 \times 10^8 \text{ K}$ ,  $kT \approx 30 \text{ keV}$   
 (n,  $\gamma$ )- ( $\gamma$ ,n) equilibrium

$$\Rightarrow N_{(A)} * \sigma_n(A) = N_{(A+1)} * \sigma_n(A+1)$$

branching points induced by  $\beta_b = f(T_S)$

$$\lambda^*/\lambda^0 > 1 \Rightarrow T_S$$



other example:  $^{205}\text{Tl}$  (stable);  $^{205}\text{Tl}^{81+} \rightarrow (\beta_b) ^{205}\text{Pb}^{81+*}$  ( $T_{1/2} = 1 \text{ y?}$ )

**key-nuclei: Supernova outbursts, galactic cosmic rays**

$^{56}\text{Ni}$ ,  $^{54}\text{Mn}$ ,  $^{44}\text{Ti}/^{44}\text{Ca}$

$^{56}\text{Ni } 0^+ \text{ (6 d)}$  \_\_\_\_\_

$^{54}\text{Mn } 3^+ \text{ (300 d)}$  \_\_\_\_\_

$Q_{\text{EC}} = 0.4 \text{ MeV}$   
 $1^+$  \_\_\_\_\_  
 $1.7 \text{ MeV}$

$Q_{\text{EC}} = 0.5 \text{ MeV}$   
 $2^+$  \_\_\_\_\_  
 $0.8 \text{ MeV}$   
 $^{54}\text{Fe } 0^+$  \_\_\_\_\_

$^{56}\text{Co } 4^+$  \_\_\_\_\_

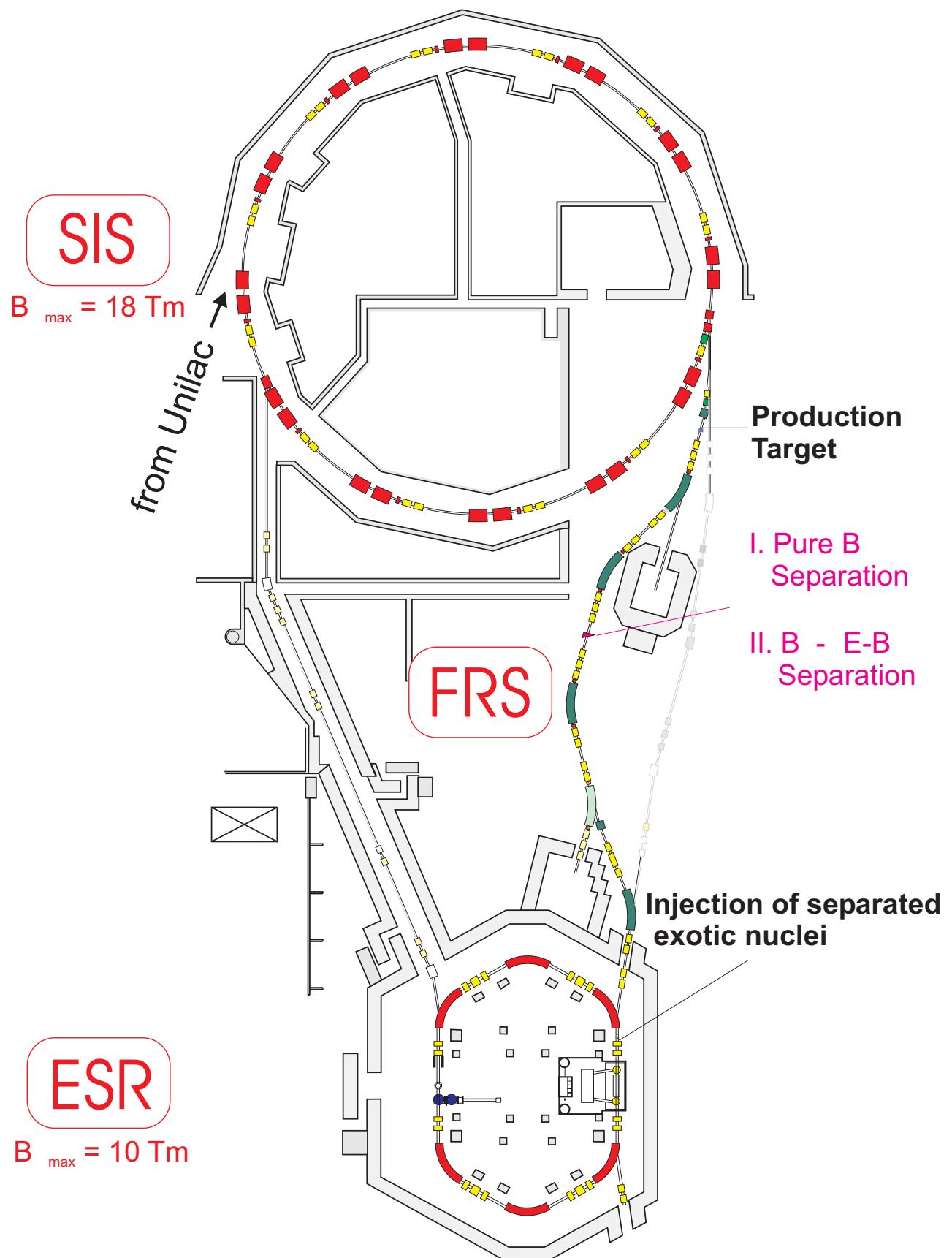
$^{54}\text{Cr } 0^+$  \_\_\_\_\_

$^{44}\text{Ti } 0^+ \text{ (60 y)}$  \_\_\_\_\_  
 $Q_{\text{EC}} = 0.01 \text{ MeV}$   
 $1^+$  \_\_\_\_\_  
 $0.15 \text{ MeV}$   
 $^{44}\text{Sc } 2^+ \text{ g.s. (4 h)}$  \_\_\_\_\_

$Q_{\text{EC}} = 2.5 \text{ MeV}$   
 $2^+$  \_\_\_\_\_  
 $E_\gamma = 1.15 \text{ MeV}$

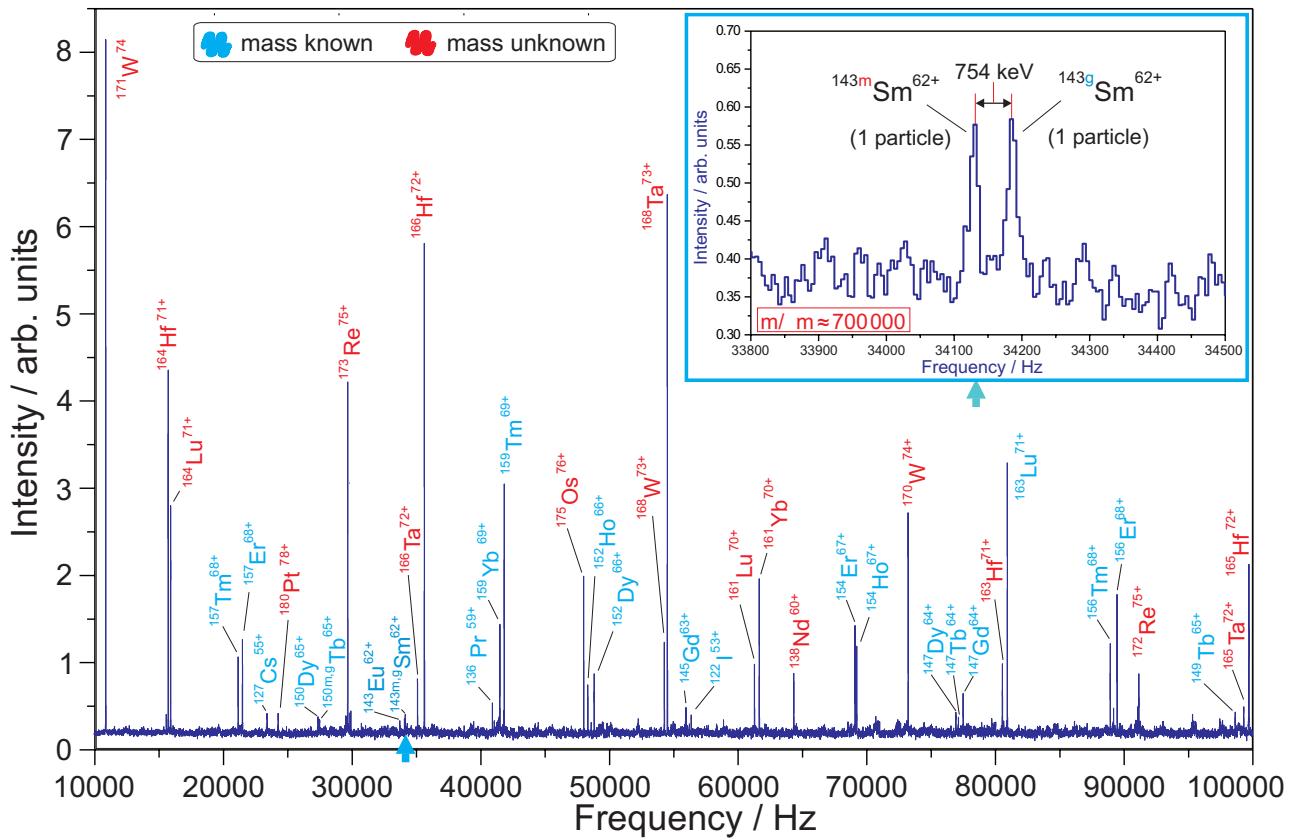
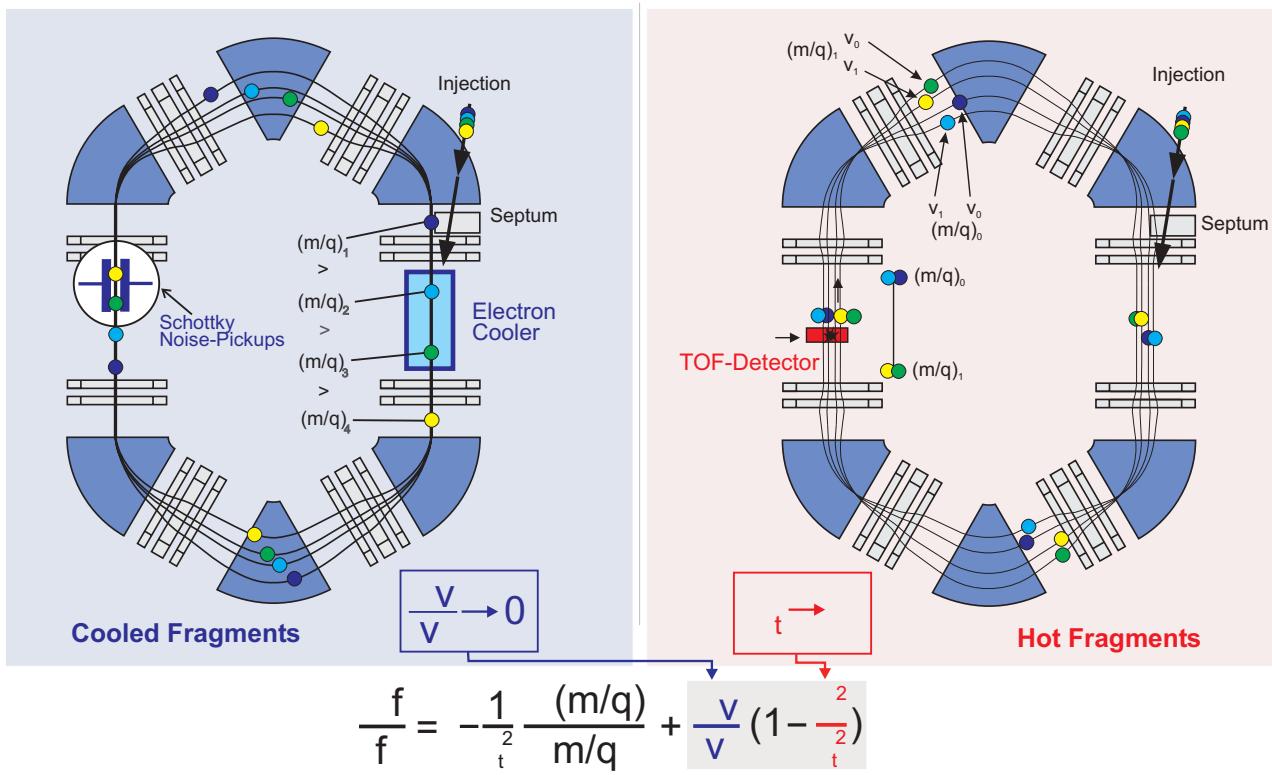
$^{44}\text{Ca } 0^+$  \_\_\_\_\_

# Mass and Lifetime Measurements of Stored Exotic Nuclei



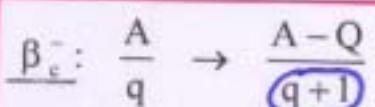
## SCHOTTKY MASS SPECTROMETRY

## ISOCHRONOUS MASS SPECTROMETRY

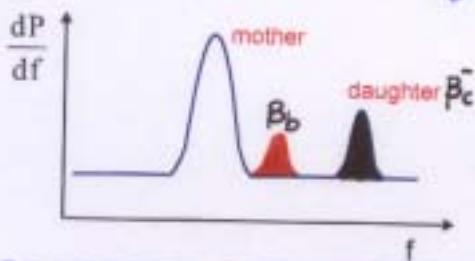


### Principle of Lifetime Measurements for Cooled Fragments

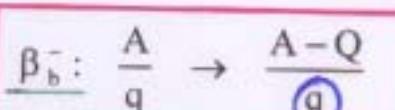
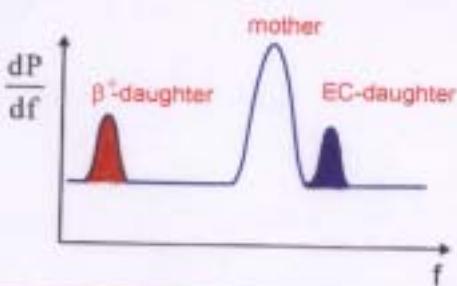
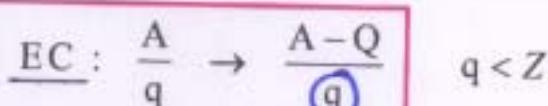
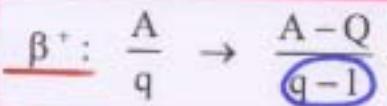
projectiles  $\longrightarrow$  FRS  $\longrightarrow$  ESR  $\longrightarrow$  Schottky Sp.



$$Q \gtrsim 1 \text{ MeV}$$

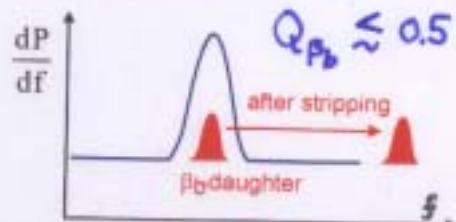


For  $Z \gtrsim 28$  both are circulating within the acceptance of the ESR

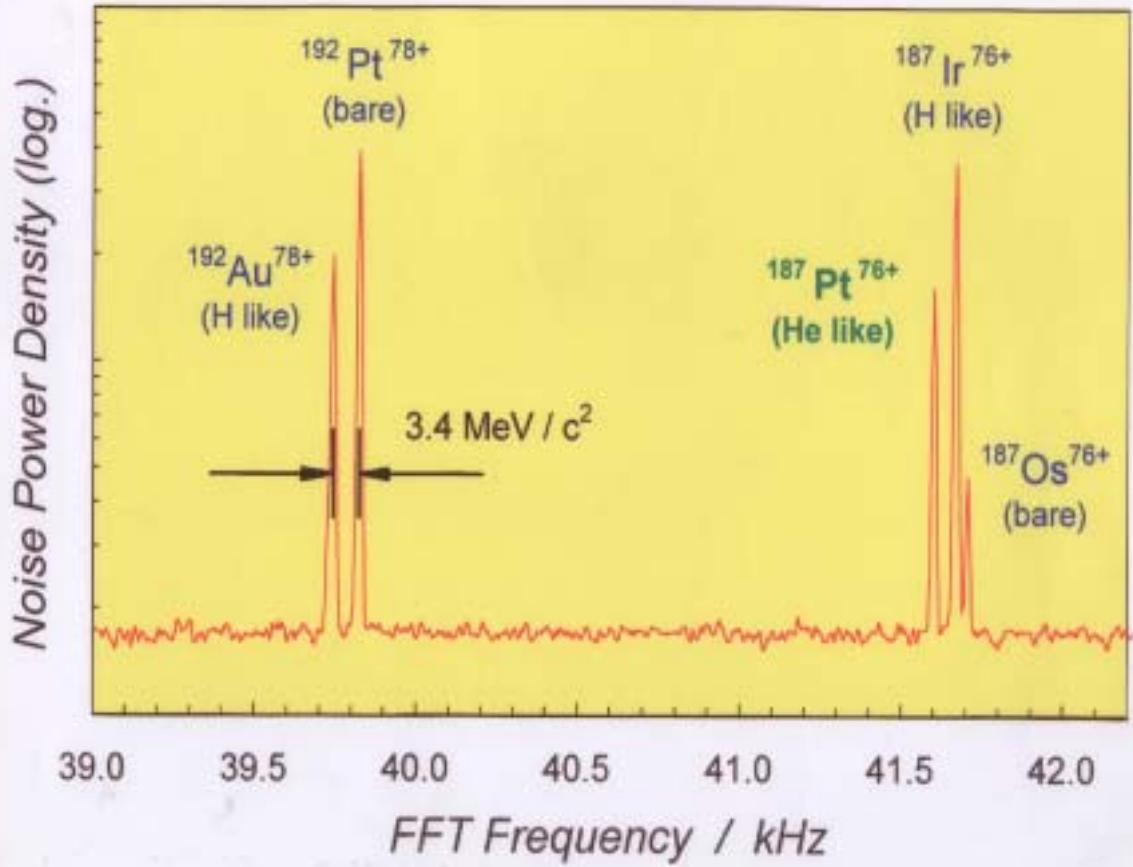


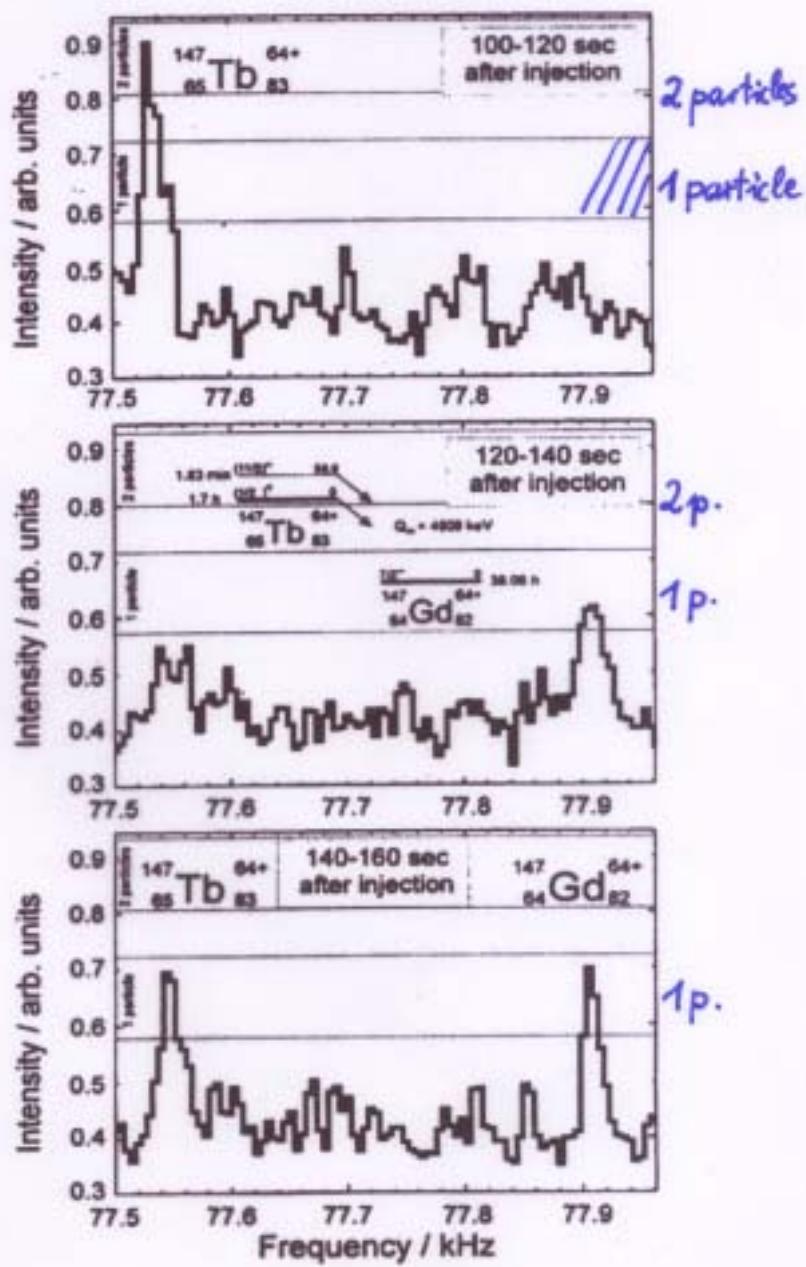
$$Q_{\beta^+} \gg Q_{\text{EC}}$$

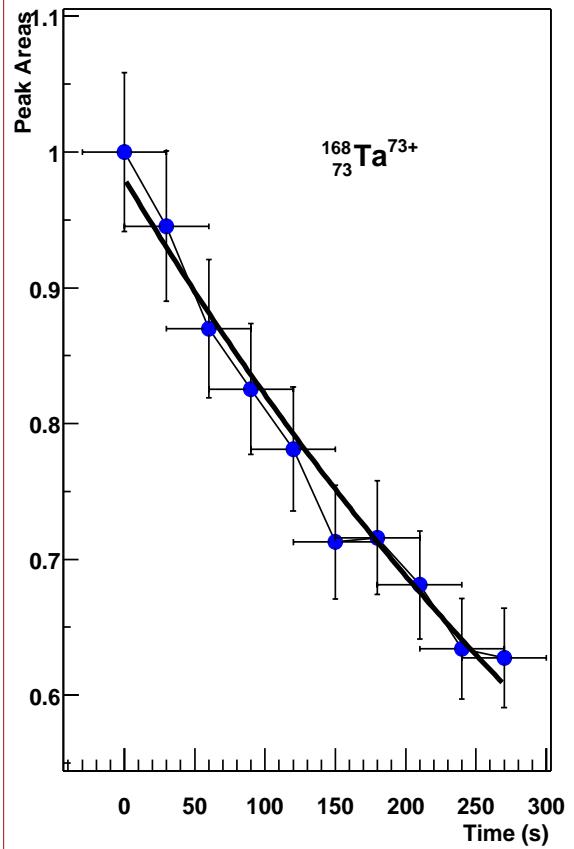
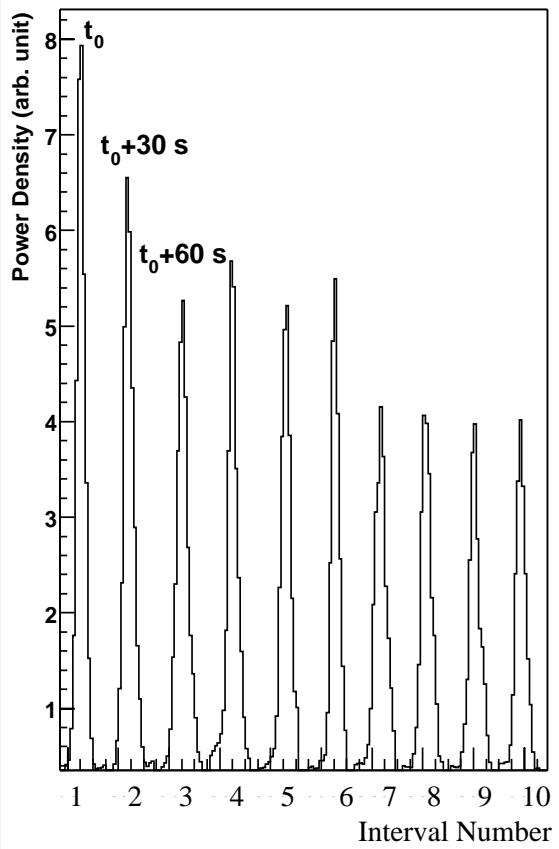
Unique:



- 1) Both mother and daughter nuclei can be measured simultaneously
- 2)  $T_{1/2}$  can be studied for different atomic charge states

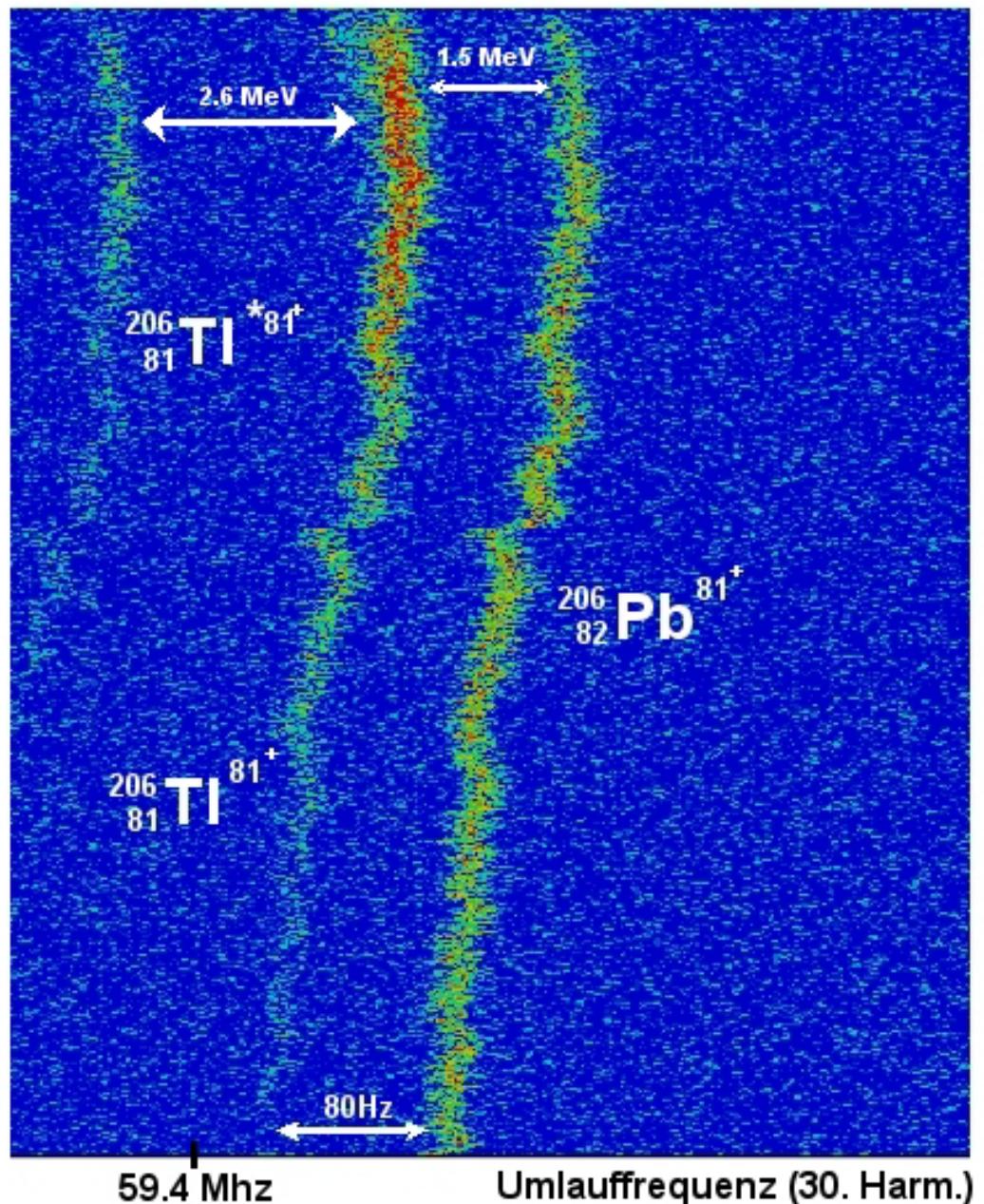


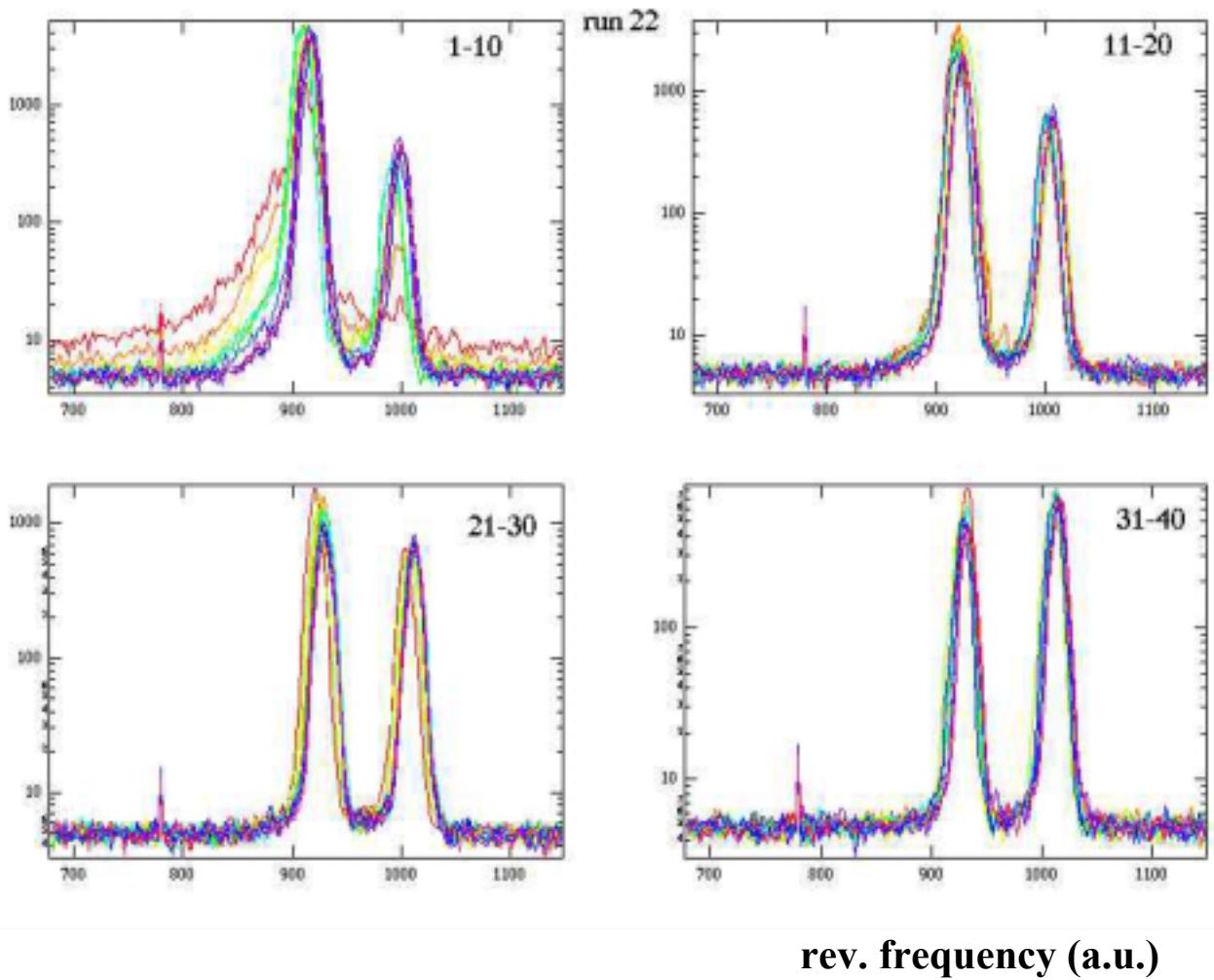




Zeit nach  
Einschuß

30 min





$$Q_{\beta b} = 1.5 \text{ MeV}$$

**1 cycle = 32.4 s (starting at injection)**

separation of  $^{207}\text{TI}/^{207}\text{Pb}$ -lines  $\approx$  **80 Hz / 60 Mhz**

**At new storage rings same techniques:**

**stochastic pre-cooling, electron cooling**

**detection of  $\beta^-$ ,  $\beta^+$ ,  $\beta_b$  and EC by**

**Schottky spectrometry / particle detectors**

**but:**

**$\approx 10^2$  more primary,  $\approx 10^4$  more secondary beams**

